# The Westerbork HI survey of spiral and irregular galaxies

# III: HI observations of early-type disk galaxies

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**Abstract.** We present H<sub>I</sub> observations of 68 early-type disk galaxies from the WHISP survey. They have morphological types between S0 and Sab and absolute B-band magnitudes between -14 and -22. These galaxies form the massive, high surface-brightness extreme of the disk galaxy population, few of which have been imaged in H<sub>I</sub> before. The H<sub>I</sub> properties of the galaxies in our sample span a large range; the average values of M<sub>HI</sub>/L<sub>B</sub> and D<sub>HI</sub>/D<sub>25</sub> are comparable to the ones found in later-type spirals, but the dispersions around the mean are larger. No significant differences are found between the S0/S0a and the Sa/Sab galaxies. Our early-type disk galaxies follow the same H<sub>I</sub> mass-diameter relation as later-type spiral galaxies, but their effective H<sub>I</sub> surface densities are slightly lower than those found in later-type systems. In some galaxies, distinct rings of H<sub>I</sub> emission coincide with regions of enhanced star formation, even though the average gas densities are far below the threshold of star formation derived by Kennicutt (1989). Apparently, additional mechanisms, as yet unknown, regulate star formation at low surface densities. Many of the galaxies in our sample have lopsided gas morphologies; in most cases this can be linked to recent or ongoing interactions or merger events. Asymmetries are rare in quiescent galaxies. Kinematic lopsidedness is rare, both in interacting and isolated systems. In the appendix, we present an atlas of the H<sub>I</sub> observations: for all galaxies we show H<sub>I</sub> surface density maps, global profiles, velocity fields and radial surface density profiles.

**Key words.** Surveys – Galaxies: fundamental parameters – Galaxies: ISM – Galaxies: kinematics and dynamics – Galaxies: spiral

# 1. Introduction

The discovery that H<sub>I</sub> rotation curves of spiral galaxies remain flat far outside the regions where the luminous mass is concentrated (Bosma 1978), has led to the now widely accepted view that galaxies are embedded in large halos of invisible matter (Faber & Gallagher 1979; Bosma 1981; van Albada & Sancisi 1986). The nature of this dark matter still remains a mystery, and poses one of the most persistent challenges to present day astronomy.

A systematic study of rotation curves in spiral galaxies, covering a large range of luminosities, morphological types and surface brightnesses, is crucial for a proper understanding of the distribution of dark matter in galaxies. Most H1 studies in recent years have, however, focused on late-type, low-luminosity systems (de Blok et al. 1996; Swaters et al. 2002); H1 observations of high-luminosity and early-type disk galaxies are rare. The only study so far aimed at a systematic investigation of H1 rotation curves over the full range of morphological types was that by Broeils (1992). However, in his sample of 23 galaxies, only one was of morphological type earlier than Sb

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and only four had  $V_{\rm max} > 250 \, {\rm km \, s^{-1}}$ . S0 and Sa galaxies were also under-represented in the study by Persic et al. (1996); their Universal Rotation Curve is based on over 1000 rotation curves of which only 2 are of type Sab or earlier.

The main reason for the lack of information on the H<sub>I</sub> properties of early-type disks is the fact that they generally have a low H<sub>I</sub> flux (Roberts & Haynes 1994). The only large-scale H<sub>I</sub> survey directed specifically at S0 and Sa galaxies was carried out by van Driel (1987), but his study was severely hampered by the low signal-to-noise ratio of his data. Jore (1997) obtained deep VLA observations for a sample of Sa galaxies, but focused mostly on asymmetries and evidence for interactions and mergers (Haynes et al. 2000).

Early-type disk galaxies form the high-mass, highsurface-brightness end of the disk galaxy population (Roberts & Haynes 1994). While the late-type spiral and irregular galaxies at the other end of the population are generally believed to be dominated by dark matter (Carignan & Freeman 1988; Broeils 1992; Swaters 1999), little yet is known about the relative importance of dark and luminous matter in earlytype disks. Knowledge of their dark matter content is a crucial step towards a proper understanding of the systematics of dark matter in galaxies.

The work presented here is part of a larger study aimed at investigating the properties of dark matter in early-type disk galaxies and its relation with the luminous components. A major part of our study consists of mapping the distribution and kinematics of neutral hydrogen in a large sample of these systems. As part of the WHISP project, we have observed 68 galaxies with morphological types between S0 and Sab. WHISP (Westerbork HI survey of spiral and irregular galaxies; Kamphuis et al. 1996; van der Hulst et al. 2001) is a survey of galaxies in the northern sky using the Westerbork Synthesis Radio Telescope (WSRT), aiming at a systematic investigation of the H<sub>I</sub> content and large-scale kinematics of disk galaxies. Observations were carried out between 1993 and 2002, and the survey now contains data for almost 400 galaxies, covering Hubble types from S0 to Im. Earlier results from WHISP were presented by Swaters et al. (2002, hereafter Paper I), Swaters & Balcells (2002, Paper II) and García-Ruiz et al. (2002).

In this paper, we present the first results of the HI observations of the early-type disk galaxies. We discuss the basic HI properties of these galaxies and we present an atlas of HI surface density maps, velocity fields, global profiles and radial surface density distributions. The data presented here will be used in a forthcoming paper to derive rotation curves, which will then be combined with optical data to study the dark matter content and distribution in these systems.

The structure of this paper is as follows. In Sect. 2 we describe the selection of the sample and some basic properties of the selected targets. In Sect. 3, the observations and data reduction steps are discussed. In Sect. 4, we describe the H<sub>I</sub> properties of our sample galaxies, and compare them to their optical properties. Notes on individual galaxies are given in Sect. 5. Section 6 presents a summary of the main conclusions. The atlas of the H<sub>I</sub> observations is presented in the appendix.

# 2. Sample selection

All galaxies presented here were selected from the WHISP survey (Kamphuis et al. 1996; van der Hulst et al. 2001). The galaxies in WHISP were selected from the Uppsala General Catalogue of Galaxies (UGC; Nilson 1973). To ensure that the galaxies would be well resolved with the WSRT, target galaxies were selected on the basis of their position on the sky ( $\delta > 20^{\circ}$ ), and angular size  $(D_{25} > 1')$ . A lower line flux limit, defined as  $f = F_{\rm HI}/W_{20}$  with  $F_{\rm HI}$  the total flux in Jy km s<sup>-1</sup> and  $W_{20}$  the width of the H<sub>I</sub> profile in km s<sup>-1</sup>, was chosen to ensure that the channel maps have a sufficient signal-to-noise ratio to produce HI surface density maps and velocity fields. For the largest part of the survey, carried out before 1999, this limit was set at f > 100 mJy. This led to the inclusion of only 11 earlytype galaxies, compared to over 300 intermediate and late-type galaxies. In 1999, a more powerful correlator and cooled frontends on all WSRT telescopes improved the sensitivity for 21cm line observations by approximately a factor of 3. For the observations with the new setup, we lowered the flux limit to f > 20 mJy.

From the final WHISP sample, we selected all 68 galaxies with morphological type between S0 and Sab. Some basic properties of these galaxies are listed in Table 1.

Our sample is not intended to be complete in any sense, but it is still interesting to compare some general properties of our objects to the full sample of galaxies in the UGC. In Fig. 1, we show in bold lines the distribution of our sample galaxies over a number of parameters. We also show the distribution of the entire UGC catalogue (thin lines), as well as all galaxies from the UGC with morphological type between S0 and Sab (dashed lines).

The histogram for morphological type simply reflects our selection criterion for early-type disk galaxies. Compared to the entire early-type disk subsample from the UGC, we have a fair sample of Sa and Sab galaxies; S0's are however underrepresented in our sample. This is due to the well-known fact that S0 galaxies have generally an even lower H1 content than Sa's (Roberts & Haynes 1994); even with the current high sensitivity of the WSRT, only a handful of S0 galaxies can be observed in single 12h runs.

Due to our diameter and H<sub>I</sub> flux limits, we select predominantly nearby galaxies, as is visible in the top-middle histogram in Fig. 1. Most of our sample galaxies have  $V_{\text{hel}} \leq 5000 \, \text{km s}^{-1}$ .

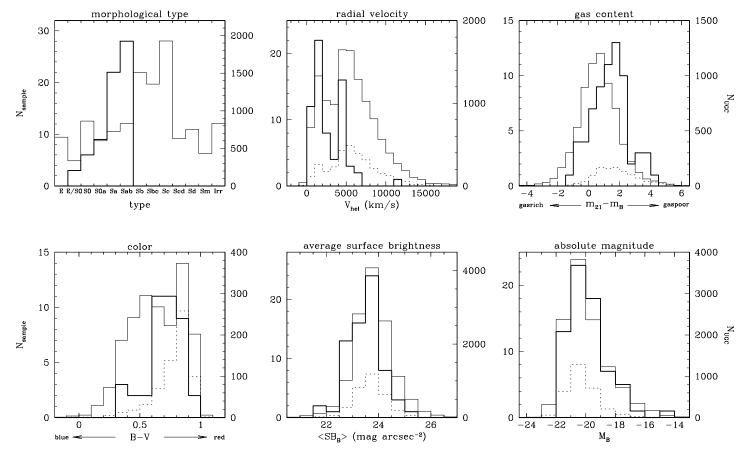
Compared to the average early-type disk galaxies in the UGC, our objects are on average slightly more gas-rich (probed by the  $m_{21} - m_{\rm B}$  'color') and marginally bluer. The distribution of average B-band surface brightness of our galaxies, defined within the 25th magnitude isophote, shows a small excess on the bright side, but this is hardly significant. Similarly, we have a small excess of low-luminosity systems.

In conclusion, our sample contains a fair number of Sa and Sab galaxies, whose properties seem to resemble those of average early-type disk galaxies in the UGC very well. Their bluer color and higher gas content may indicate a higher level of star formation, but the differences for both parameters are small. The fact that the average surface brightness and luminosity distributions are almost identical to those for the entire subsample of early-type disks from the UGC, confirms our conjecture that we are not looking at a special subclass of these systems. The situation for S0's is less clear, due to the small number contained in this sample.

# 3. Observations and data reduction

All observations described here were carried out with the Westerbork Synthesis Radio Telescope (WSRT) in single 12 hour sessions. 11 galaxies were observed before 1999, using the old front-ends and correlator. The remaining 57 were observed after the major upgrade of the WSRT, using new cooled front-ends on all telescopes and a new broadband correlator. A summary of the observational parameters is given in Table 2.

For most observations, the field of view was centered on the target galaxy. In a few cases, more than one galaxy was observed in one pointing, and some galaxies lie off-center in the field of view (notably UGC 94, 499, 508, 6621, 8699 and 12815). The angular resolution depends on the coverage of the UV-plane during the observations. Furthermore, as the WSRT is an east-west array, the resolution in declination depends on



**Fig. 1.** Characteristics of our sample galaxies (bold lines), compared with the entire UGC (thin lines) and with all galaxies from the UGC with type between S0 and Sab (dashed lines). The scales on the left of each panel give the distribution of the galaxies in the sample presented here, the scales on the right apply to the galaxies from the UGC. All data are taken from LEDA<sup>1</sup>, except the heliocentric velocities for our sample galaxies, which are taken from this study. Note that some parameters (e.g. color, gas content) are not available for all galaxies in LEDA; the total number of galaxies included is different for each histogram.

the position on the sky and scales as  $1/\sin \delta$ . The resulting beam size for each galaxy is given in Table 2.

Observations were usually done with 128 channels; the total bandwidth was chosen on the basis of the H<sub>I</sub> line width  $W_{20}$ , such that we had optimal wavelength resolution while still having a reasonable baseline to construct continuum maps. The resulting velocity resolution, before Hanning smoothing, is thus either  $\sim 2.5$ , 5, 10 or 20 km s<sup>-1</sup>, with the exact value depending on the central wavelength of the band (see Table 2).

The data reduction consisted of several stages; each will be described below.

#### 3.1. standard WHISP pipeline

The initial reduction steps were carried out following the standard WHISP data reduction pipeline. The raw UV data were inspected and calibrated using the NEWSTAR software package. Bad data-points were flagged interactively. The UV data were then Fourier transformed to the image plane, and antenna patterns were calculated. For each set of observations, data cubes were created at 3 different resolutions and corresponding noise levels. One is at full resolution using all available

UV data-points; the resulting beam size is given in Table 2. The other two were produced by down-weighting progressively more long baselines from the UV data. While doing so, the spatial resolution was decreased to respectively  $\approx 30^{\prime\prime}$  and  $\approx 60^{\prime\prime}$ , but the signal-to-noise level for extended emission was strongly enhanced.

All further data reduction was done with GIPSY (Groningen Image Processing System; Vogelaar & Terlouw 2001, http://www.astro.rug.nl/~gipsy/). For all data cubes, the continuum was subtracted by fitting, at each line of sight, a first order polynomial to the channels without HI emission. The cube at 60" resolution was then Hanning smoothed in velocity and masks were created by hand to identify the regions containing HI emission. The data cubes were then CLEANed (Schwarz 1978), using the masks to define the search areas. CLEANing was iterated down to 0.5 times the rms in each channel map. The rms noise in the cleaned, full resolution channel maps,  $\sigma_{\rm ch}$ , is given in Table 2.

In more than 50% of the cases, the data cubes at full resolution are of sufficient quality to derive useful H<sub>I</sub> maps and velocity fields. However, in cases where the signal is very weak, the signal-to-noise ratio in the full resolution data may be too low to extract the emission line profiles. In those cases we used

<sup>&</sup>lt;sup>1</sup> LEDA: Lyon Extragalactic Database, http://leda.univ-lyon1.fr/

the smoothed data in the subsequent analysis. Additionally, the smoothed data cubes are sometimes used to study the distribution and kinematics of the gas in the faint outer regions. In the figures in the atlas we indicate which data cube was used for each galaxy.

# 3.2. defining masks

Before proceeding to the next steps, we defined first the regions that contain emission; parts of the data cubes that contain only noise were masked out, such that as little noise as possible entered the global profiles and surface density maps.

Creating the masks is not straightforward and requires special care. Applying a simple sigma-clipping criterion on the channel maps is not sufficient, as we may miss extended low-level emission. Therefore, we first smoothed the channel maps: the channel maps at full, 30 and 60" resolution were smoothed to 30, 60 and 120" respectively. Subsequently the data cubes were Hanning smoothed in velocity. The signal-to-noise ratio for extended emission in the resulting smoothed data cubes is much higher and low-level emission is now easily detected.

The Hanning-smoothed data cubes at 30, 60 and 120" were then used to create masks for the original channel maps at full, 30 and 60" resolution respectively; the masks were created in 2 steps. First we selected all pixels in the smoothed channel maps that have intensity  $> 2\sigma$ . In this way, we did not only select real emission from the galaxy, but also many noise peaks in the data. In the second step we selected by hand the emission which we deemed to be real. Emission was defined to be real if it is spatially extended (i.e.  $\geq 2$  beam areas) and present in more than 2 adjacent channels or when it is very bright ( $\geq 4\sigma$ ). Small regions that are only slightly brighter than  $2\sigma$  are defined to be noise and were rejected. This second step is by definition subjective. Some peaks we judged to be noise may actually be real and vice versa. However, since we are working on extensively smoothed data, the real emission is more easily distinguished from noise, and the effect is small.

The resulting clipped channel maps were then used as masks for the original data cubes.

#### 3.3. global profiles and total HI masses

Global profiles were derived from the data cubes at 60" resolution. In these data cubes, we are most sensitive to extended H<sub>I</sub> emission, and the flux enclosed by the masks is maximal. The global profiles were constructed by adding up all flux in each masked channel map and correcting for primary beam attenuation. Care was taken not to include flux from companion galaxies.

The error in each point of the profile is related to the size of the masked region in the channel map. Due to the non-uniform sampling of the UV-plane by the WSRT, the noise in the channel maps is not spatially independent and the error does *not* simply increase as the square root of the mask area. Instead, we established an empirical relation between mask size and error in the flux and used the result to estimate the errors on our global profiles.

The global profiles are shown in the top right panels in the atlas (Appendix A).

Line widths were determined at the 20% and 50% levels. If a profile was double peaked, the peaks on both sides were used separately to determine the 20% and 50% levels, otherwise the overall peak flux was used. The profile widths  $W_{20}$  and  $W_{50}$  were defined as the difference between the velocities at the appropriate level on each side of the profile. We followed Verheijen & Sancisi (2001) to correct the observed line widths for instrumental broadening, assuming an internal velocity dispersion of the gas of  $10 \text{ km s}^{-1}$ :

$$W_{20}^{c} = W_{20} - 35.8 \left[ \sqrt{1 + \left(\frac{R}{23.5}\right)^{2}} - 1 \right]$$

$$W_{50}^{c} = W_{50} - 23.5 \left[ \sqrt{1 + \left(\frac{R}{23.5}\right)^{2}} - 1 \right],$$
(1)

with R the instrumental velocity resolution in km s<sup>-1</sup>.

The systemic velocities  $V_{\rm sys}$  were defined as the average of the velocities at the profile edges at the 20% and 50% level. Distances D were derived using a Hubble constant of 75 km s<sup>-1</sup> Mpc<sup>-1</sup>, correcting the systemic velocities for Virgocentric inflow using the values given by LEDA. The profile widths, systemic velocities and distances are given in Table 3.

The total H<sub>I</sub> mass was derived from the total flux, i.e. the integral of the global profile just derived, and is given by:

$$M_{\rm HI} = 236 \cdot D^2 \cdot \int F \, \mathrm{dv},\tag{2}$$

where  $M_{\rm HI}$  is in solar masses, D is the distance in Mpc, F is the primary beam corrected flux in mJy, and the integral is over the total bandwidth. The derived total fluxes and H<sub>I</sub> masses are given in Table 3 as well.

# 3.4. integrated H<sub>I</sub> maps

HI maps were created by integrating the masked data cubes along the velocity direction and correcting the result for primary beam attenuation. The noise in these maps is not constant: the galaxies presented here are generally characterized by very large velocity gradients in the center and slowly varying velocities in the outer parts, so that we are summing over a variable number of channel maps at different positions. For each map, we estimated an average noise level,  $\sigma_{\text{map}}$ , following the prescriptions in Verheijen & Sancisi (2001). Noise fields were constructed based on the number of channel maps contributing to the integrated H<sub>I</sub> maps at each position and the rms noise in the individual channel maps. Using these fields, we selected all points in our H<sub>I</sub> maps with signal-to-noise ratio between 2.75 and 3.25. The average value of these points was then divided by 3 to obtain the  $\sigma_{\rm map}$  value. It is important to note that this is an average noise level; low-level emission which is only present in a few adjacent channels may be significant, even if it is weaker than the  $2\sigma_{\text{map}}$ -level of the outer contour in the figures.

The HI column density maps are shown in the left hand panels in the atlas. In Table 2, we give the corresponding values of  $\sigma_{\rm map}$  in atoms cm<sup>-2</sup>.

#### 3.5. velocity fields

The two methods most commonly used to derive velocity fields from HI data cubes are to calculate the intensity weighted mean (IWM) velocity of the line profiles or to fit Gaussians to them. For the data presented here, neither of these methods proved to be adequate in recovering the true radial velocity of the gas. Particularly in the central regions of the earlytype galaxies studied here, large velocity gradients are present and the line profiles are seriously skewed by beam-smearing. In those cases, both the IWM velocity and the central velocity of the fitted Gaussians can be offset significantly from the true gas velocity at the projected radius (e.g. Paper I). Here, we derive velocity fields by fitting a Gauss-Hermite polynomial that includes an h3-term for skewness to the line profiles (van der Marel & Franx 1993). The resulting velocities from this method lie usually closer to the peak in the profile than simple IWM or Gauss fits, especially in the central regions.

We created velocity fields in two steps. In the first step, we fitted the Gauss-Hermite function to the profiles in the masked data cubes. However, since the masked data cubes contain no noise information outside the emission lines, it is difficult to determine the significance and reliability of the fitted profiles. To overcome these difficulties, we used the fitted parameters from the first step as initial estimates for a second fit to the Hanningsmoothed, unmasked line profiles. For these latter fits we could develop strict criteria to determine whether or not a profile represents true emission: only those fits were accepted that have 1) a line strength greater than 2.5 times the rms noise  $\sigma_h$  in the Hanning-smoothed channel maps and 2) a velocity within the range defined by the global profile. These two constraints were usually sufficient to exclude erroneous results. As an extra sanity check, we excluded fits with 3) a profile width smaller than 4 or larger than 200 km s<sup>-1</sup>. This criterion was, however, rarely

It is important to note here that the sigma-clipping criterion in I) above was performed using the rms in the Hanningsmoothed channel maps directly. This rms is not the same as the  $\sigma_{\rm map}$  level in the HI surface density maps, which was determined empirically (see previous section). Thus, it is possible that points which fall outside the  $2\sigma_{\rm map}$  contour in the HI density map, have a fitted line strength greater than  $2.5\,\sigma_{\rm h}$  and vice versa.

The final velocity fields are shown in the top middle panels in the atlas. In the bottom middle panels, we show the fitted velocities over-plotted on position-velocity slices along the major axes through the data cubes. From these figures, it is clear that the inclusion of an h3-term for skewness in the fitting procedure worked well in recovering the true velocity of the gas for most galaxies (e.g. UGC 89, 6786). In some cases, however, in particular in edge-on or poorly resolved galaxies, projection effects lead to such complicated line-profiles that even skewed Gaussians fail to recover the projected rotational velocity of the gas. In our sample, the velocity fields of UGC 1310, 2045, 5906, 5960, 6001, 6742, 7704, 8271 and 12713 suffer seriously from this effect and can therefore not be used to derive rotation curves; for UGC 2183, 3354, 4605 and 11670, only the inner regions of the velocity fields are affected.

# 3.6. radial profiles, HI diameters and surface densities

The integrated H<sub>I</sub> maps created in Sect. 3.4 were used to derive radial profiles of the H<sub>I</sub> surface density. The intensities were azimuthally averaged in concentric elliptical annuli. Pixels without measured signal in the total H<sub>I</sub> map were counted as zero. The azimuthal averaging was also done separately for the approaching and receding half of the galaxy to obtain a crude estimate of the level of asymmetries present in the H<sub>I</sub> map.

The orientation of the annuli was determined with different methods, depending on the galaxy. For a number of galaxies, tilted ring fits to the velocity field were available. The details of these fits will be described in a forthcoming paper dealing with the rotation curves of our galaxies. If no tilted ring analysis was available (for example if the galaxy is too poorly resolved, or if the HI gas does not show signs of circular rotation), we used the ellipticity and position angle from LEDA. However, in some cases LEDA does not give a position angle or is clearly inconsistent with our data. Especially for near face-on galaxies, LEDA's position angles – if present – can be far off from the angles we infer from the velocity field. In those cases, we did not use LEDA's position angle, but rather determined it by eye from our HI maps and velocity fields.

The measured intensities were corrected to face-on surface densities, and converted to units of  $M_{\odot}$  pc<sup>-2</sup>. The resulting radial profiles are shown in the bottom right panels in the atlas.

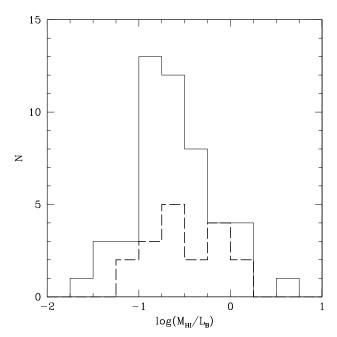
For edge-on galaxies, this method does not work, as emission from different radii is superposed onto the same position on the sky. For these cases, we used the iterative Lucy deconvolution method developed by Warmels (1988b). The HI surface density maps were integrated parallel to the minor axis to get HI strip integrals. These were then converted to surface density profiles using the iterative deconvolution scheme from Lucy (1974), under the assumption of axisymmetry for the gas distribution.

From the radial profiles, H<sub>I</sub> radii were determined. The H<sub>I</sub> radius  $R_{HI}$  is defined as the radius where the face-on corrected surface density drops below a value of  $1~M_{\odot}~pc^{-2}$  (equivalent to  $1.25 \cdot 10^{20}~atoms~cm^{-2}$ ). The derived radii are given in Table 3.

Finally, we derived the average H<sub>I</sub> surface density within the H<sub>I</sub> radius ( $<\Sigma_{\rm HI}>_{R_{\rm HI}}$ ), as well as within the optical radius ( $<\Sigma_{\rm HI}>_{R_{25}}$ ). The optical radii were derived from the absorption corrected 25th B-band mag arcsec<sup>-2</sup> diameters, taken from LEDA. The average H<sub>I</sub> surface densities are also listed in Table 3.

# 4. HI properties of early-type disk galaxies

As mentioned in the introduction, the data presented here are intended for a study of rotation curves and dark matter in early-type disk galaxies. In this paper, however, we focus on the global H<sub>I</sub> properties of our sample galaxies, and the relation with their optical characteristics; several aspects of this analysis are discussed below.



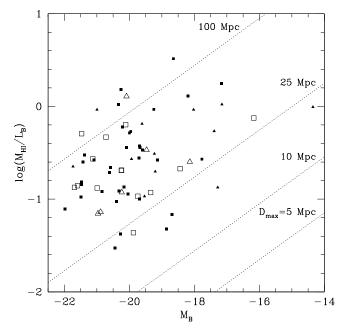
**Fig. 2.** The distribution of  $M_{HI}/L_B$ . The solid line indicates all galaxies of type Sa and Sab, the dashed line indicates the S0's.

#### 4.1. Hi content

In Fig. 2, we show the distribution of  $M_{\rm HI}/L_B$ . For  $L_B$ , we used the absolute B-band magnitudes from LEDA (column (6) in Table 1) and a solar absolute magnitude in the B-band of 5.47 (Cox 2000). For comparison, we have split the sample in S0/S0a's and Sa/Sab's.

The average values of  $\log(M_{\rm HI}/L_B)$  for Sa/Sab's and S0/S0a's is  $-0.62 \pm 0.44$  and  $-0.50 \pm 0.40$  in solar units respectively, where the errors give the standard deviations of the respective distributions. These values are higher than the average values for early-type disks found by Roberts & Haynes (1994), confirming that we have predominantly selected gasrich galaxies in our sample (cf. Sect. 2 and Fig. 1). There is, however, a large variation in the HI content of our galaxies; the most gas-rich galaxies contain about 100 times more gas, compared to their stellar luminosity, than the most gas-poor systems. This range is larger than seen in most late-type galaxies, which span at most only one order of magnitude in gas content (Roberts & Haynes 1994; Broeils & Rhee 1997; Paper I).

The large spread in HI content is also apparent in Fig. 3, where we plot the HI mass-to-light ratio versus B-band luminosity. In later-type galaxies, a general trend exists, such that the more luminous galaxies contain relatively less HI (Verheijen & Sancisi 2001; Paper I). In our sample, this correlation is virtually absent. There seems to be a lack of low-luminosity galaxies ( $M_B \gtrsim -18.5$ ) with low relative gas content, but this is probably a selection effect. The dotted lines in Fig. 3 give the maximum distance  $D_{max}$  out to which galaxies would be included in our sample.  $D_{max}$  is defined as the distance where the flux density of a galaxy, given by  $f = F_{HI}/W$ , drops below 20 mJy. W, the profile width, is estimated from the absolute magnitude using the Tully-Fisher relation found by Verheijen (2001) and assuming an average inclination of



**Fig. 3.** M<sub>HI</sub>/L<sub>B</sub> versus absolute B-band magnitude. Sa/Sab galaxies are indicated by squares, S0/S0a's by triangles. The filled symbols indicate galaxies whose morphology is not or only mildly asymmetric, open symbols denote moderately or strongly lopsided galaxies. The dashed lines indicate the maximum distance at which a galaxy would still be included in our sample (see text).

 $60^{\circ}$ . It is obvious that high-luminosity galaxies with a low gas content can be observed out to much larger distances than low-luminosity systems. Low-luminosity ( $M_B \gtrsim -18.5$ ) galaxies with low relative gas content ( $\log(M_{\rm HI}/L_B) \lesssim -1$ ) may be as common as their high-luminosity counterparts, the lack of low-luminosity, low  $M_{\rm HI}/L_B$  galaxies simply being the result of the smaller volume that our survey covers for these systems.

The large spread in gas content could be explained if a substantial fraction of our early-type disk galaxies have recently accreted gas from small dwarf galaxies. The spread in gas content would then reflect the accretion history of these galaxies. Indeed, many of our galaxies show signs of recent interaction or accretion (see Table 3), but the degree of asymmetry in the gas distribution does not correlate with the gas fraction, as can be seen in Fig. 3. It seems unlikely that recent accretion or merger events alone can explain the observed spread in gas content of these galaxies.

Another possible explanation for the large spread lies in the fact that H<sub>I</sub> content is a property of the disk, whereas our galaxies contain an additional bulge component. If indeed bulges contribute to the optical luminosity but are not related to the H<sub>I</sub> content, variations in bulge-to-disk luminosity ratios would lead to additional scatter in plots like Figs. 2 and 3. We have checked this hypothesis for a number of galaxies in our sample for which we have accurate optical photometry and bulge-disk decompositions at our disposal, but found that the scatter in gas content did not decrease significantly when, instead of the total luminosity, we only considered the luminosity of the disk component.

Thus, the HI content in these early-type disk galaxies seems to be rather independent of any basic optical property. In Sect. 4.3, we will show that the HI content is correlated with the relative extent of the gas disks, in the sense that the gas disks of gas-rich galaxies are more extended, but of similar surface density, than those of gas-poor systems. Although this implies that the large spread in gas content is, at least partly, related to the large range in the relative sizes of the gas disks (cf. Sect. 4.2), it still offers no explanation as to why some galaxies have larger gas disks, and correspondingly larger HI masses, than others. The origin of the large spread in gas content in these systems remains therefore puzzling.

# 4.2. Hi diameters

In Fig. 4, we show the distribution of  $D_{HI}/D_{25}^{\textit{B,c}}$ , with  $D_{HI}$  the diameters of the Hı disks, measured at the 1  $M_{\odot}$  pc $^{-2}$  level.  $D_{25}^{\textit{B,c}}$  are the optical diameters, measured at the absorption corrected 25th B-band mag arcsec $^{-2}$  level, taken from LEDA. The average ratio of Hı to optical diameter for the 49 Sa/Sab galaxies in our sample is  $1.72\pm0.70$ ; as above, the error gives the standard deviation of the distribution. This average value is identical, but with larger spread, to the value found by Broeils & Rhee (1997) for a sample of 108 spiral galaxies, mostly of intermediate and late type.

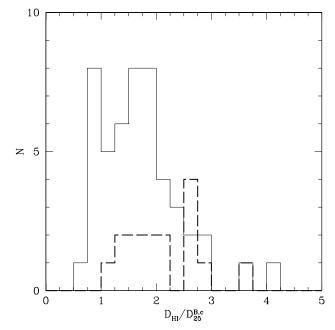
9 Sa/Sab galaxies (18%) in our sample have  $D_{\rm HI} < D_{25}^{B,c}$ . Most of these cases show clear signs of interaction (e.g. UGC 4862, 5559, 6621) and it seems likely that these galaxies have lost part of their gas disks in the encounters. This would be consistent with the results of Cayatte et al. (1994), who showed that galaxies in the center of the Virgo cluster have smaller HI disks than those in the field and argued that this is caused by ram-pressure stripping and interactions between the cluster members. However, some galaxies with small gas disks in our sample seem undisturbed (e.g. UGC 7489 or 11914); the origin of the small extent of their HI disks is unclear.

The average ratio of H<sub>I</sub> to optical diameter for the S0/S0a's in our sample is  $2.11 \pm 0.70$ . This is slightly higher than for the Sa/Sab's, but the number of objects is too small to be able to say whether this is significant. Note that three S0/S0a galaxies in our sample (UGC 2154, 3426 and 4637) have an azimuthally averaged H<sub>I</sub> surface density less than  $1~M_{\odot}~pc^{-2}$  everywhere, such that their H<sub>I</sub> radius is not defined; the gas in these three galaxies is clearly distorted and probably originates from a recent merger or interaction event.

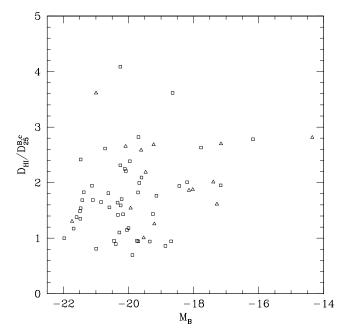
The ratio  $D_{\rm HI}/D_{25}^{B,c}$  between H<sub>I</sub> and optical diameters shows no clear correlation with optical luminosity (Fig. 5). The lack of low-luminosity galaxies with small H<sub>I</sub> disks is probably the result of the same selection effects as described in the previous section.

#### 4.3. Hi surface brightness

Previous studies have revealed a tight relation between the total H<sub>I</sub> mass and H<sub>I</sub> diameter of a galaxy, with  $M_{\rm HI} \propto (D_{\rm HI})^{\alpha},~\alpha \approx 1.9$  (Broeils & Rhee 1997; Verheijen & Sancisi 2001; Paper I). The small scatter around this relation, and the



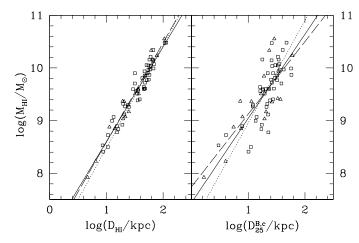
**Fig. 4.** The distribution of  $D_{\rm HI}/D_{25}^{B,c}$ . The solid line indicates all galaxies with type Sa and Sab, the dashed line indicates the S0's.



**Fig. 5.** Ratio of H<sub>I</sub> to optical diameter  $D_{HI}/D_{25}^{B,c}$  versus absolute B-band magnitude. Sa/Sab galaxies are indicated by squares, S0/S0a's by triangles.

fact that the slope is so close to 2, imply that the average H<sub>I</sub> surface density within the 1  $M_{\odot}$  pc<sup>-2</sup> isophote is almost universal from galaxy to galaxy. The correlation of H<sub>I</sub> mass with optical diameter is well defined too, but with larger scatter.

These relations are also present in the early-type galaxies studied here, as can be seen in Fig. 6. The relation between  $M_{\rm HI}$  and  $D_{\rm HI}$  is indistinguishable from the one found in previous



**Fig. 6.** Total H<sub>I</sub> mass versus H<sub>I</sub> (*left*) and optical (*right*) diameter. Sa/Sab galaxies are indicated by squares, S0/S0a's by triangles; the solid and long-dashed lines give the corresponding fits. The dotted lines give the relations found by Broeils & Rhee (1997).

studies (the dotted line shows the relation from Broeils & Rhee (1997)). We find:

$$\begin{array}{ll} log(M_{HI}/M_{\odot}) = 1.80 \, log(D_{HI}/kpc) + 6.80 & (Sa/Sab) \\ log(M_{HI}/M_{\odot}) = 1.89 \, log(D_{HI}/kpc) + 6.70 & (S0/S0a). \end{array} \eqno(3)$$

The scatter around each of the fits is 0.15 and 0.14 dex respectively, comparable to the dispersion found by Broeils & Rhee (1997).

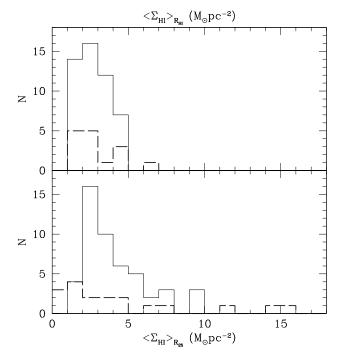
We find a slightly shallower relation between H<sub>I</sub> mass and optical diameter:

$$\begin{split} \log(M_{\rm HI}/M_{\odot}) &= 1.60 \log(D_{25}^{\it B,c}/\rm kpc) + 7.44 & (Sa/Sab) \\ \log(M_{\rm HI}/M_{\odot}) &= 1.38 \log(D_{25}^{\it B,c}/\rm kpc) + 7.74 & (S0/S0a). \end{split} \tag{4}$$

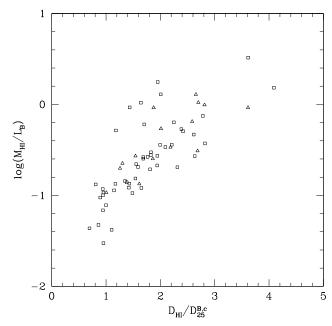
The shallow slopes in these relations could again be the result of a selection bias in our survey. Small, low H1-mass galaxies are unlikely to be included in our sample, which explains why the galaxies in our sample with  $D_{25}^{B,c} < 10 \, \mathrm{kpc}$  lie systematically above the relation found by Broeils & Rhee (1997). Given our selection criteria, we see no reason to conclude that early-type disk galaxies follow a different  $M_{\mathrm{HI}} - D_{25}^{B,c}$  relation than later-type spirals. Note however that the scatter, 0.34 and 0.38 dex for the Sa/Sab and S0/S0a samples respectively, is larger than the one found by Broeils & Rhee (1997).

In Fig. 7, we show the distributions of average H<sub>I</sub> surface density within the H<sub>I</sub> and optical diameters. The small scatter in the  $M_{\rm HI}-D_{\rm HI}$  relation translates into a narrow distribution in  $<\Sigma_{\rm HI}>_{R_{\rm HI}}$ . The average values of  $<\Sigma_{\rm HI}>_{R_{\rm HI}}$  are  $2.8\pm0.8$  and  $2.9\pm1.4\,M_{\odot}\,{\rm pc}^{-2}$  for Sa/Sab and S0/S0a galaxies respectively; the errors give the standard deviations of the distributions. Thus, early-type disk galaxies have on average somewhat lower H<sub>I</sub> surface densities than later-type spirals, but with similar scatter (Broeils & Rhee 1997). The dispersion in the distribution of  $<\Sigma_{\rm HI}>_{R_{25}}$  is much larger, but the average values are again lower than in later-type spirals (cf. Roberts & Haynes 1994; Broeils & Rhee 1997).

As a last illustration that the average H<sub>I</sub> surface density shows little variation from galaxy to galaxy, we show in Fig. 8



**Fig. 7.** The distribution of average H<sub>I</sub> surface density within the 1  $M_{\odot}$  pc<sup>-2</sup> isophote (*top*) and within the optical diameter (*bottom*). The solid line indicates all galaxies with type Sa and Sab, the dashed line indicates the S0's.



**Fig. 8.** Relative gas content,  $M_{\rm HI}/L_B$ , versus relative gas diameter,  $D_{\rm HI}/D_{25}^{B,c}$ . Sa/Sab galaxies are indicated by squares, S0/S0a's by triangles.

the relation between relative gas content and the Hi-to-optical diameter ratio. This figure shows that if a galaxy has a relatively high gas content, it is because its gas disk is relatively extended; thus the average Hi surface density will still be comparable to that in more gas-poor systems. It is, however, still unclear why some galaxies have much larger Hi disks than others (i.e. are much more gas-rich, cf. Sects. 4.1 and 4.2).

# 4.4. holes, rings and the relation with star formation

Most early-type disk galaxies presented here whose H<sub>I</sub> is concentrated in a regular disk, have a central hole or depression in their H<sub>I</sub> distribution. Galaxies where we do not observe a central hole are either not well resolved, or do not have a regularly rotating disk; it is likely therefore, that *all* early-type galaxies with a regular gas disk have a central depression in their H<sub>I</sub> distribution.

In some galaxies, most of the gas is concentrated in rings. In most cases, the rings follow the orientation of the optical disks, but there are also cases (e.g. UGC 12276) where the orientation of the ring is different from that inferred from the optical image or the velocity field. Galaxies with gas rings offer an interesting possibility to study the relation between neutral gas and the stellar light distribution. In Fig. 9 we compare the radial gas distribution for 6 galaxies with distinct gas rings with their B-band luminosity profile. The gas surface density profiles are derived from the HI profiles shown in Appendix A, multiplied by 1.44 to account for the presence of other elements (mostly helium). The photometric profiles are taken from a parallel study of the stellar mass distribution in a subsample of the galaxies presented here (Noordermeer et al., in prep.).

UGC 1541 and 3546 (top panels) are barred galaxies with distinct spiral arms. In both cases there is a slight overdensity of blue light at the same radius as the gas ring; this may simply reflect the fact that gas and stars respond in the same way to the gravitational perturbations from the bar and spiral arms.

The other four galaxies, however, are not barred, and have very little spiral structure. Yet in two cases, UGC 3993 and 11914 (middle panels), the HI surface density peaks correspond to small but distinct overdensities of light at the same radii. In UGC 11914 this excess light is related to a prominent ring of blue stars in the original images (see also Buta et al. 1995). H $\alpha$  emission is detected in this ring as well (Pogge 1989), indicating that the gas is actively forming stars (Battinelli et al. 2000). We are not aware of H $\alpha$  observations of UGC 3993, but the stars in its ring are significantly bluer than in the central parts, indicating that a population of young stars may exist in this ring as well. In contrast, in UGC 2487 and 6787 (bottom panels), no excess light over the regular exponential disks is detected and the colors of the stars do not vary over the ring. Pogge & Eskridge (1993) found no H $\alpha$  emission in UGC 2487.

In an attempt to understand why there is no current star formation in the rings of UGC 2487 and 6787, while there are strong indications that it *is* happening in UGC 3993 and 11914 where the total gas densities are similar, we compared the observed gas densities to the threshold density for star formation derived by Kennicutt (1989). Assuming that star formation in spiral galaxies is regulated by the onset of gravitational instabilities in their gas disks, Kennicutt applied the criterion for disk stability from Toomre (1964) to these systems and derived the following threshold density for star formation:

$$\Sigma_{\rm crit} = \alpha \frac{\kappa c}{3.36G}.\tag{5}$$

 $\kappa$  is the epicyclic frequency, which we derived from rotation curves from tilted ring fits to the velocity fields. For c, the velocity dispersion of the gas, we followed Kennicutt, and as-

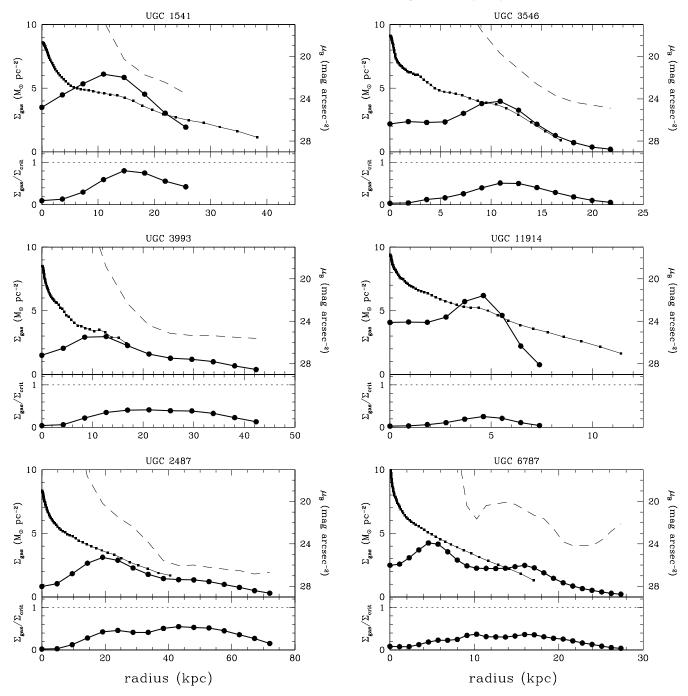
sumed a constant value of 6 km s<sup>-1</sup>. For the high surfacebrightness galaxies studied here, this value is probably on the low side, and the actual values of  $\Sigma_{crit}$  may be higher than what we derive here. However, higher values for the threshold density will only aggravate the discrepancy described below, and we choose  $c = 6 \text{ km s}^{-1}$  to be able to compare our results directly with Kennicutt.  $\alpha$  is a dimensionless quantity of order unity; from an empirical study of star formation cutoffs in spiral galaxies, Kennicutt derived a value of  $\alpha = 0.67$ . The applicability of this recipe was recently confirmed by Martin & Kennicutt (2001) for a sample of 32 nearby spiral galaxies. The critical densities in our galaxies are plotted, as a function of radius, with the dashed lines in Fig. 9. Due to the large rotation velocities and the correspondingly large  $\kappa$ , the critical densities are high and the observed gas densities are far below the threshold in all cases.

Thus, if the description of star formation thresholds by Kennicutt (1989) is correct, it is not surprising that we do not see a correlation between the H<sub>I</sub> rings and the stellar light distribution in UGC 2487 and 6787; the gas densities are simply much too low to sustain large-scale star formation. But then it is unclear why UGC 3993 and 11914 do form stars. In fact, the galaxy out of these 6 which has the highest star formation activity, UGC 11914, has the *lowest* ratio (< 0.25 everywhere) of observed to critical density!

Note that we have ignored the contribution of molecular gas to the total gas densities. It seems, however, unlikely that the presence of molecular gas can explain the discrepancy between observed and critical gas densities in these galaxies. Several CO-surveys have been carried out in recent years, and all find that molecular gas is concentrated in the inner parts of galaxies (Sofue et al. 1995; Wong & Blitz 2002; Sofue et al. 2003), with exponential scale lengths comparable to those of the underlying stellar disks (Regan et al. 2001). No substantial column densities of CO emission are generally observed at large radii. Thus, while we may have underestimated the total gas densities in the inner regions of the galaxies in Fig. 9, the effect will be small in the H<sub>I</sub> rings of interest, which all lie at large radii, and the total gas densities are truly smaller than the critical density.

In recent years, more cases have been reported of star formation in regions where the gas density is lower than the critical density from equation 5. Kennicutt (1989) and Martin & Kennicutt (2001) already noted that a fraction of the galaxies in their sample (notably M33 and NGC 2043) showed widespread star formation, even though their gas surface densities were below the threshold. van Zee et al. (1997) observed 11 dwarf galaxies and showed that, although the gas densities were below the threshold in all cases, some of them had a substantial level of star formation. Several recent papers report the detection of star formation in the outskirts of nearby galaxies, where the observed gas densities are well below the threshold densities for star formation too (Ferguson et al. 1998; Cuillandre et al. 2001; Thilker et al. 2005). Thus, a number of cases have now been identified where the simple criterion of Kennicutt (1989) does not apply.

Recently, Schaye (2004) argued that, rather than being regulated by global, gravitational instabilities, star formation requires a phase transition from warm to cold gas, and is thus



**Fig. 9.** Comparison between gas and stellar light distributions in galaxies with pronounced gas rings. UGC 1541 and 3546 are barred galaxies; the others have no bars. *Top panels*: Filled circles and bold lines show the HI surface density profiles, multiplied by 1.44 to account for the presence of helium. Radii are converted to kpc. Filled squares and thin lines show B-band photometric profiles. Dashed lines show the critical density for star formation, according to Kennicutt (1989). For UGC 11914, the critical density is larger than  $10 \, \text{M}_{\odot} \, \text{pc}^{-2}$  everywhere. *Bottom panels*: Ratio between observed gas density and critical density. The dotted line indicates the threshold for star formation.

governed by thermal instabilities instead. He predicted that the critical density for star formation is more or less independent of global properties such as rotation velocity, and has a universal value in the range  $3-10~M_{\odot}~pc^{-2}$ . Although it is not clear how this prediction compares to the large spread in observed threshold densities from Kennicutt (1989), it is suggestive that all galaxies in Fig. 9 have gas densities around the lower end of

this range, and would thus be at the verge of instability for star formation.

It seems, however, premature to completely abandon Kennicutt's theory, until more knowledge is obtained about the small-scale distribution of the gas (cf. van Zee et al. 1997). Braun (1997) studied the structure of HI disks in 11 well-resolved, nearby galaxies and found that the HI was distributed

in small, high-density regions with covering factors between 6 and 50%. Thus, if UGC 3993 and 11914 are extreme cases where the covering factor is low, and UGC 2487 and 6787 have a smoother HI distribution, the observed (lack of) correlation between HI distribution and star formation in these galaxies might still be consistent with the predictions from Kennicutt (1989). Note, however, that Kennicutt did not explicitly account for the 'patchiness' of gas in his galaxies; an average amount of substructure in the gas distribution must implicitly be included in his parameter  $\alpha$ . There is no a priori reason to assume that galaxies such as UGC 3993 or 11914 have much lower covering factors than average, but we cannot rule out this possibility, and high-resolution observations are needed to check this.

# 4.5. lopsidedness

Asymmetries in the morphology and kinematics of galaxies are still a subject of debate. Rix & Zaritsky (1995) found that about one third of all spiral galaxies show large-scale asymmetries in their optical images. Global Hi-profiles have asymmetric shapes for at least 50% of all disk galaxies (Baldwin et al. 1980; Richter & Sancisi 1994; Haynes et al. 1998), but without a full 2D mapping of the H<sub>I</sub> component in these galaxies, it is difficult to determine the origin of the asymmetries in the profiles. With the advent of a large number of spatially resolved HI images and velocity fields, it was found that not only the spatial distribution of the gas is often asymmetric, but that many galaxies are also lopsided in their kinematics: Swaters et al. (1999) estimated that about half of all disk galaxies have asymmetric kinematics. Lopsidedness seems particularly common in latetype spirals; Matthews et al. (1998) found that out of a sample of 30 extremely late-type galaxies, about 75% had clearly lopsided H<sub>I</sub> profiles.

In some cases, lopsidedness can be linked to ongoing interactions with companion galaxies. However, lopsidedness is also common in isolated galaxies (Wilcots & Prescott 2004), implying that lopsidedness may be an intrinsic feature of disk galaxies and that the underlying dark matter distribution may be asymmetric too (Jog 1997; Noordermeer et al. 2001). Even if lopsidedness in galaxies is triggered by tidal interactions or mergers, it must be a long-lived phenomenon to explain the relatively high frequency of lopsidedness in isolated systems.

A full, quantitative analysis of the occurrence of lopsidedness in the galaxies in our sample is beyond the scope of the present paper. Instead, we have determined by eye how many of our early-type galaxies are lopsided, and briefly discuss the results here. The degrees of asymmetry in the global profile, morphological appearance and kinematical structure (velocity field and major axis xv-diagram) are indicated with stars in Table 3, where 0, 1, 2 or 3 stars mean not, mildly, moderately or severely lopsided respectively. In some cases the gas is too strongly distorted or too poorly resolved to determine whether it is asymmetric. These cases are indicated with n.a. (not available).

Asymmetries in the global profiles and gas distribution are very common in the early-type disk galaxies studied here.

Respectively 35 and 34 galaxies (51 and 50%) have at least mildly asymmetric global profiles or surface density maps. In most cases, however, the asymmetries can be related to the presence of or interactions with companion galaxies. *All* galaxies, except one (UGC 5060), which have severely lopsided gas distributions, and most cases where the gas distribution is moderately lopsided, show clear signs of ongoing interactions. On the other hand, morphological lopsidedness is rare in early-type disk galaxies without nearby companions or signs of recent interactions. Of all 41 galaxies that do not show clear signs of interaction, only 15 (37%) are morphologically asymmetric, usually only mildly. Only 3 isolated galaxies (7%) are more than mildly asymmetric.

Kinematical lopsidedness is rare in early-type galaxies in general, interacting or not. Out of all the galaxies for which we can determine the degree of symmetry in the kinematics (53), UGC 624 is the only one with a severely lopsided velocity field; in this case the lopsidedness can be explained as a result of tidal interaction with its nearby companion UGC 623. The only two cases of isolated galaxies with significantly lopsided kinematics, UGC 5960 and UGC 11852, are both extreme galaxies in our sample in the sense that the former is one of the least luminous galaxies in our sample ( $M_B = -17.39$ ), while the latter has the most extended gas disk, relative to the optical diameter, of all our sample galaxies ( $D_{\rm HI}/D_{25}^{B,c} = 4.1$ ).

We conclude that lopsidedness in the HI distribution is at least as common in early-type disk galaxies as in later-type galaxies; truly symmetric gas disks are rare. In most cases, however, morphological asymmetries are clearly related to ongoing interactions or accretion. Morphological lopsidedness in isolated early-type disks, and kinematical asymmetries, are much rarer, supporting the findings of Matthews et al. (1998).

# 5. notes on individual galaxies

Many galaxies in our sample show interesting features in their H<sub>I</sub> distribution or kinematics which deserve special emphasis. We discuss these below.

UGC 89 (NGC 23) and UGC 94 (NGC 26) are separated by only 10'. They form a loose group together with Scd galaxy UGC 79 (20' to the southwest) and 4 other galaxies (Garcia 1993). A few small H<sub>I</sub> clouds are visible in the neighbourhood. Optical counterparts are visible on the DSS for most of them, so they are probably dwarf companions of the main group members. A tidal tail-like structure is visible south-east of UGC 94, indicating a recent interaction or merger event. In our low-resolution data cube, this tidal feature is seen to be connected to the main disk of UGC 94 and extends about 10' (~ 180 kpc) to the east.

UGC 499 (NGC 262), also known as Markarian 348, is a well-known Seyfert 2 galaxy (Koski 1978). We detect H<sub>I</sub> emission out to distances of 100 kpc from the center, but it is clearly distorted and does not rotate regularly around the center. UGC 508 (NGC 266) lies about 23′ to the northeast of UGC 499; the peculiar structure of the gas in UGC 499 might be a result of a past interaction with its neighbour. The H<sub>I</sub> properties of UGC 499 were discussed in detail by Heckman et al. (1982)

and Simkin et al. (1987).

**UGC 624** (NGC 338) is one of the most lopsided galaxies in our sample, both in the gas distribution and in the kinematics. The asymmetry could be a result of tidal interaction with nearby neighbour UGC 623. Note also the small 'blob' of gas just southeast of the main gas disk.

UGC 1310 (NGC 694) resides in a small group of galaxies, the main other members of which are NGC 680, 691 and 697 (Garcia 1993). The central regions of this group were imaged in H<sub>I</sub> before by van Moorsel (1988). Apart from UGC 1310, we also detect H<sub>I</sub> emission in NGC 691, 697 and UGC 1313, but as these are all of later morphological type, they are not included in the present study. The H<sub>I</sub> disk of UGC 1310 is barely resolved in our observations. As a result, the line profiles are strongly affected by beam-smearing; the method described in Sect. 3.5 to derive the velocity field is therefore inadequate and fails to recover the projected rotational velocities of the gas.

Most of the gas in **UGC 2045** (NGC 972) seems to be concentrated in a regularly rotating disk that has the same orientation as dustlanes in the optical image. The line-profiles in this disk have complicated shapes and the method described in Sect. 3.5 for the derivation of the velocity field fails to recover the projected rotational velocities of the gas. In the outer parts, several large filaments of gas are detected; the gas in these filaments clearly does not follow the rotation of the inner parts.

UGC 2154 (NGC 1023) is the brightest galaxy of a nearby group of 13 galaxies (Tully 1980). The total H<sub>I</sub>-flux for this galaxy is quite large (80.13 Jy km s<sup>-1</sup>), but the gas is scattered over a large area, has a very low column-density and does not reside in a regular disk. It seems most likely the result of a recent merger or accretion event, possibly with one of the other group members (cf. Sancisi et al. 1984).

Most of the gas in UGC 2183 (NGC 1056) is concentrated in an edge-on disk which coincides with the dust lane seen in the optical image. Due to this edge-on orientation, projection effects lead to strongly non-Gaussian line profiles; the method described in Sect. 3.5 to derive the velocity field is therefore inadequate and fails to recover the projected rotational velocities of the gas in the central parts. At larger radii, the gas seems to warp out of the plane and the orientation becomes more face-on

**UGC 2487** (NGC 1167) is a nice example of an S0 galaxy with an extended and regularly rotating H<sub>I</sub> disk. The gas seems to be in circular motion at radii up till 80 kpc. The central hole does not reflect a true absence of gas, but is rather the result of absorption against a central continuum source.

The companion east of **UGC 2916** is PGC 14370. When the data are smoothed to lower resolution, some emission is seen to bridge the space between the two galaxies, indicating that some interaction is going on. There are clear signs of off-planar gas in PGC 14370, but the main disk of **UGC 2916** seems relatively undisturbed.

**UGC 2953** (IC 356) is by far the best resolved galaxy in our sample. The gas is concentrated in pronounced spiral arms that extend far beyond the bright optical disk. The velocity field shows distinctive 'wiggles' in the isovelocity contours where

the gas crosses the arms.

The H<sub>I</sub> distribution, velocity field and global profile of UGC 3205 are almost perfectly symmetric, except for the twisting of the isovelocity contours in the bar region.

UGC 3354 is almost perfectly edge-on and has a peculiar double warp. Both the stellar and gas disks seem to warp 'clockwise' first. At larger radii, where no starlight can be detected, the gas disk reverses its warp to the opposite direction. Projection effects lead to complicated line-profiles in the inner parts of the disk; the method described in Sect. 3.5 to derive the velocity field is inadequate there and fails to recover the projected rotational velocities of the gas.

The gas distribution in **UGC 3426** is irregular and seems to be connected to the gas disk of UGC 3422,  $\approx 100$  kpc to the north-west. The projected surface density of the gas is very low, with a maximum of only about  $0.3 \, M_{\odot} \, pc^{-2}$ . It appears as if the gas in **UGC 3426** has been tidally drawn out of the gas disk of UGC 3422.

UGC 3642 is a very peculiar galaxy. The gas inside the bright optical disk seems to be regularly rotating in the same plane as defined by the stellar light distribution. But further out, the direction of motion is reversed. Each component individually appears undisturbed and the gas seems to be in regular rotation in both the inner and the outer regions. Two options seem possible to explain the peculiar kinematics of this galaxy. The first is that both components are fully decoupled, in which case this galaxy would carry some resemblance with the 'Evil Eye' galaxy (NGC 4826), which also has two decoupled gas disks (Braun et al. 1992; Rubin 1994). However, both the total HI mass and the size of the disks in UGC 3642 are an order of magnitude larger than in NGC 4826. Alternatively, the gas disk could be extremely warped, such that the direction of motion along the line of sight is reversed. We are unable to distinguish between the two options, and more detailed observations and modelling are needed to clarify the nature of this galaxy. Note that the total extent of the outer gas disk is extremely large, with  $D_{\rm HI} > 100$  kpc.

At first sight, **UGC 3965** (IC 2204) seems a similar case as UGC 3642, with the kinematical position angle changing by almost 180°. But in this case, the observed velocity field can also be explained by an almost face-on disk with a mild warp. While the gas would be moving in the same direction everywhere, its radial velocity along the line of sight could easily be reversed.

UGC 4605 (NGC 2654) is almost perfectly edge-on in the inner parts, but the gas seems mildly warped in the outer parts. The position-velocity diagram shows that the rotation velocities decline strongly towards the edge of the HI disk; the decline starts already before the onset of the warp and must reflect a true decrease in the rotation velocities. In the very inner parts, the line-profiles are strongly affected by projection effects and the method described in Sect. 3.5 to derive the velocity field fails to recover the projected rotational velocities of the gas there.

**UGC 4637** (NGC 2655) is a nearby Seyfert 2 galaxy whose nuclear regions show complex structure, both in high-resolution

radio continuum observations as in optical emission lines (Keel & Hummel 1988). These authors attribute the complexity of the central parts to a recent interaction or merger event. This interpretation seems to be confirmed by the large-scale structure of the neutral gas seen in our observations. The HI is clearly disturbed, with a large extension to the northwest; the gas does seem to have a general sense of rotation, but it is clearly not on regular circular orbits. The optical image of this galaxy shows some distinct loops and shells, further confirming the hypothesis of a recent interaction.

**UGC 4666** (NGC 2685) is also known as the Spindle or Helix galaxy, because of its prominent polar ring. Early H<sub>I</sub> observations of this galaxy were presented by Shane (1980). Our observations, at higher sensitivity and resolution, confirm both the outer gas ring as well as H<sub>I</sub> related to the helix-like structure inside this ring. The gas in the helix is kinematically distinct from the rest of the galaxy and seen at almost right angles to the outer ring.

UGC 4862 (NGC 2782) has a large, banana-shaped tail of gas northwest of the main disk. The gas in the central regions is clearly disturbed as well and does not show signs of regular rotation around the center. The optical image is peculiar too, with some striking shells. These facts were already noted by Smith (1994), who interpreted the peculiar structure as evidence for a recent merger of the main galaxy with a low-mass companion. Schiminovich et al. (1994, 1995) studied HI in shells around elliptical galaxies. It seems not unreasonable to assume that UGC 4862 is a similar case as their galaxies, but in an earlier stage of its evolution.

The main disk of **UGC 5253** (NGC 2985) is regular in its distribution and kinematics. In the outer parts, a large one-armed spiral of gas is seen to extend towards a small 'blob' of gas southwest of the galaxy. **UGC 5253** lies about 20' west of the late-type spiral UGC 5316, which also has a disturbed morphology (e.g. Vorontsov-Vel'Yaminov 1977). The redshifts of the two systems differ by about 250 km s<sup>-1</sup> only, so the peculiar structures of both galaxies could well be the result of a tidal interaction.

UGC 5559 (NGC 3190) is part of Hickson compact group of galaxies 44, together with UGC 5554, 5556 and 5562 (Hickson et al. 1989). All group members lie in our field of view, but the observations for these galaxies were done with the 'old' Westerbork receivers and the sensitivity is insufficient to detect any emission in UGC 5554 and 5562. UGC 5559 itself is only barely detected, and the signal-to-noise ratio of the data is too low to make a detailed study of the gas distribution in this galaxy. Only for the more gas-rich, late-type spiral UGC 5556 useful Hr surface density maps and velocity fields could be obtained, but as this galaxy does not meet our selection criteria, it is not discussed here further.

**UGC 5906** (NGC 3380) is poorly resolved and the line-profiles suffer from beam-smearing. As a result, the method described in Sect. 3.5 to derive the velocity field is inadequate and fails to recover the projected rotational velocities of the gas.

**UGC 5960** (NGC 3413) is close to edge-on and projection effects lead to non-Gaussian line profiles. As a result, the method

described in Sect. 3.5 to derive the velocity field is inadequate and fails to recover the projected rotational velocities of the gas.

Due to the poor resolution and resulting beam-smearing in our observations of **UGC 6001** (NGC 3442), the method described in Sect. 3.5 to derive the velocity field is inadequate and fails to recover the projected rotational velocities of the gas.

The central hole in the disk of **UGC 6118** (NGC 3504) is not due to a true absence of gas, but rather an artefact caused by H<sub>I</sub> absorption against a central continuum source (see also UGC 2487).

UGC 6621 (NGC 3786) and UGC 6623 (NGC 3788) are clearly interacting. A giant tidal tail of H<sub>I</sub> extends almost 50 kpc northward from UGC 6623. On our optical image, a lowsurface-brightness counterpart is visible. Although the gas of both galaxies seems to overlap in the space between them, it is separated in velocity and we can distinguish which gas belongs to which galaxy in the full resolution data cube. Thus we were able to generate radial profiles for the gas in each galaxy separately. At 60" resolution, the data become too heavily smoothed and the regions of emission from the individual galaxies start overlapping in individual channel maps as well. We cannot make global profiles of the emission of each galaxy separately anymore at this resolution, and we were forced to generate the global profiles from the 30" data. But at 30" already the galaxies are hardly resolved and the masks for this resolution do not miss much flux compared to those of the 60" data.

**UGC 6742** (NGC 3870) is poorly resolved and the line-profiles suffer from beam-smearing. As a result, the method described in Sect. 3.5 to derive the velocity field is inadequate and fails to recover the projected rotational velocities of the gas.

The gas in the inner parts of **UGC 6786** (NGC 3900) follows the light distribution of the optical image. Most of the gas in the outer parts is concentrated in two diffuse spiral arms which seem to be warped with respect to the inner parts.

The well-studied galaxy **UGC 7166** (NGC 4151) is one of the original Seyfert galaxies (Seyfert 1943). The H<sub>I</sub> in this galaxy follows the stellar light in the bar and the spiral arms (cf. Pedlar et al. 1992).

**UGC 7256** (NGC 4203) has a peculiar, filamentary gas distribution. The gas has a general sense of motion around the center, but it is clearly not on regular circular orbits. This galaxy bears some resemblance with UGC 4637 (see above).

**UGC 7489** (NGC 4369) has the smallest HI disk, relative to its optical diameter, of all galaxies in our sample  $(D_{\rm HI}/D_{25}^{B,c}=0.86)$ , with the exception of a few interacting systems or galaxies where the HI or optical radii are not well defined (e.g. UGC 5559 or 10448). In the optical image, very weak spiral structure and dust absorption can be seen at the locations of the HI emission.

**UGC 7704** (NGC 4509) has a peculiar morphology. The optical image has a position angle of about 155°, but the H<sub>I</sub> map seems to be almost perpendicular to this, with a position angle of about 57°. The kinematical major axis is different again,

about 37°. **UGC 7704** does not seem to have any major companions which could cause this strange structure, so the true nature of it remains unclear. Due to the poor resolution, the line-profiles in this galaxy suffer from beam-smearing and the method described in Sect. 3.5 to derive the velocity field fails to recover the projected rotational velocities of the gas.

No HI gas is detected in the bulge of **UGC 7989** (NGC 4725), nor in the giant bar. It is rather concentrated in narrow spiral arms, which coincide with the stellar arms. The outer arm on the east side is particularly prominent, whereas the arms on the northwest side seem to be truncated, causing marked asymmetries in the HI surface density map and global profile.

Most of the gas in UGC 8271 (NGC 5014) seems to be concentrated in a ring which is tilted with respect to the main stellar disk. Close inspection of the optical image reveals a faint polar ring-like structure, and the H<sub>I</sub> at the corresponding locations seems to be rotating in the same plane as defined by this ring. A tail of gas extends further to the south. 12' to the northnortheast, a number of large HI clouds are detected, all without optical counterparts. UGC 8271 resides in a loose group, the main other members of which are NGC 5005 and 5033. The peculiar phenomena in the gas distribution and kinematics in and around UGC 8271 may all be related and suggest a recent merger or accretion event, possibly with a small gas-rich member of the same group of galaxies. Due to the complicated structure of the gas in this galaxy and the resulting projection effects, the line-profiles are strongly non-Gaussian and the method described in Sect. 3.5 to derive the velocity field fails to recover the projected velocities of the gas.

UGC 8805 (NGC 5347) is at first sight a similar case as UGC 7704. The position angle of the HI-disk seems to be almost perpendicular to that of the optical disk. Note though that we only get sufficient signal at 60" resolution; the optical disk then fits almost entirely in one resolution element. There is a hint that the HI is actually elongated along the optical image in the central regions, but observations at higher resolution and sensitivity are required to confirm this. If indeed the gas is aligned with the stars in the inner parts, UGC 8805 could be an example of a galaxy with an extreme warp. This would also explain the strong decline of radial velocities along the major axis, as shown in the xv-slice.

The optical image of **UGC 8863** (NGC 5377) is completely dominated by a giant bar. Two faint spiral arms are seen to extend from its edges. The H<sub>I</sub> emission nicely follows the bar and spiral arms; the latter are much more pronounced here. In the center, bar-induced streaming motions produce a twist in the isovelocity contours.

**UGC 9133** (NGC 5533) has a large one-armed spiral in the outer regions which seems to be warped with respect to the main disk. The inner parts are highly regular. The xv-diagram indicates that the rotation velocities are declining. In the outer parts, this can be explained as a result of the warp, but the decline sets in well within the radius of the warp and must reflect a truly falling rotation curve.

The H<sub>I</sub> which we detect in **UGC 10448** (NGC 6186) is very peculiar. The gas seems offset from the optical image by

about 30". Close inspection of the optical image reveals a faint galaxy behind the main disk. Thus, we are seeing here the coincidental alignment of two spiral galaxies. This is further confirmed by the redshift of the HI emission of about 11350 km s<sup>-1</sup>, whereas the optical redshift of the foreground galaxy is 2935 km s<sup>-1</sup> (UZC; Falco et al. 1999). The foreground galaxy has been detected at the correct redshift in HI as well (Rosenberg & Schneider 2000), but the UGC contains the redshift of the background system and confused us into observing at the wrong frequency. As no morphological classification, nor photometric data, are available for the background galaxy, we have excluded this galaxy from the statistical analysis in Sect. 4, and merely show the data for general interest.

UGC 11269 (NGC 6667) must recently have undergone a merger. A giant one-armed spiral arm seems to extend out from the main disk. It is partly visible in the optical image as well. Most of this material seems to be virialized already, since it follows closely the velocities of the gas in the inner parts. Note also the peculiar, asymmetric shape of the position-velocity diagram along the major axis: at the receding side, the peak velocity is much higher than at the approaching side. On both sides, the rotation velocities decline strongly away from the center

Most of the gas in **UGC 11670** (NGC 7013) is concentrated in the bar and the spiral arms. In the outer regions, the gas distribution is irregular, with several filaments emanating from the main disk. The line profiles in the inner parts are strongly non-Gaussian; the method described in Sect. 3.5 to derive the velocity field is inadequate there and fails to recover the projected rotational velocities of the gas.

The gas in **UGC 11914** (NGC 7217) is predominantly concentrated in a ring. The stellar population in the ring is distinctively bluer than the surroundings, indicating that the gas is associated with an enhanced level of star formation (see Sect. 4.4).

Most of the gas in **UGC 12276** (NGC 7440) is concentrated in a ring which is elongated perpendicular to the optical and kinematical major axis. It seems unlikely that this peculiar geometry can be explained solely by the presence of the little companion  $\approx 40$  kpc to the east, so the true nature of it remains unclear.

Due to the poor resolution and resulting beam-smearing in the observations of **UGC 12713**, the method described in Sect. 3.5 to derive the velocity field is inadequate and fails to recover the projected rotational velocities of the gas.

**UGC 12815** (NGC 7771) is strongly distorted and seems to be interacting with its neighbours UGC 12813, NGC 7771A and UGC 12808. Two giant tidal tails extend to more than 100 kpc away from the system.

# 6. Conclusions

In the previous sections, we have presented the results of a study of the H<sub>I</sub> properties of a sample of 68 early-type disk galaxies, with morphological type ranging from S0 to Sab and absolute B-band magnitude between -14 and -22. A number of conclusions can be drawn:

- The HI content of early-type disk galaxies is highly variable. There is a wide range in HI mass-to-light ratios M<sub>HI</sub>/L<sub>B</sub>, with the most HI rich galaxies in our sample containing about 2 orders of magnitude more gas, relative to the stellar luminosity, than the most gas-poor systems. The average values for log(M<sub>HI</sub>/L<sub>B</sub>) are −0.62 ± 0.44 and −0.50 ± 0.40 in solar units for Sa/Sab and S0/S0a galaxies respectively. The errors give the standard deviations of the distributions.
- The average ratio  $D_{HI}/D_{25}^{B,c}$  between H1 and optical diameter, defined at  $1~M_{\odot}~pc^{-2}$  and 25 B-band mag arcsec<sup>-2</sup> respectively, is  $1.72\pm0.70$  for Sa/Sab galaxies and  $2.11\pm0.70$  for S0/S0a's. These values are comparable to those observed for later-type spiral galaxies, but with larger spread.
- The average H<sub>I</sub> surface brightness in our sample galaxies is slightly lower than that in later-type galaxies. Within our sample, the variations from galaxy to galaxy are small. The average values for  $<\Sigma_{\rm HI}>_{R_{\rm HI}}$ , the effective H<sub>I</sub> surface brightness inside the 1  $M_{\odot}$  pc<sup>-2</sup> isophote, are 2.8 ± 0.8 and 2.9 ± 1.4 for Sa/Sab and S0/S0a galaxies respectively.
- All early-type galaxies whose gas is distributed in a regular rotating disk have a central hole or depression in their HI distribution.
- A number of galaxies in our sample have distinct, axisymmetric rings of gas. In some of these cases, the overdensity of gas coincides with regions of enhanced star formation and a population of young stars, even though the gas densities are far below the threshold for star formation derived by Kennicutt (1989). In other cases no star formation activity is present, even though the gas densities are comparable. These discrepancies suggest the existence of an additional regulation mechanism for star formation at low gas densities, the exact nature of which still needs to be clarified.
- Morphological and kinematical peculiarities are very common in early-type disk galaxies, often related to ongoing or recent interaction events. In many galaxies, we see indications for a tidal origin of the gas. Interactions with neighbour galaxies seem to be a driving force in the evolution of the neutral gas component in many early-type disk galaxies.
- Many early-type disk galaxies have lopsided morphologies.
   In most cases the asymmetries can be explained as the result of interaction, accretion or merger events. Few isolated galaxies have lopsided gas distributions.
- Kinematic lopsidedness is rare in early-type disk galaxies, even in interacting systems.

The data presented in the atlas in the appendix form the basis for a study of the rotation curves and dark matter content in these early-type disk galaxies. This study is currently ongoing, and the results will be presented in a forthcoming paper.

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**Table 1.** WHISP early-type disk galaxy sample: basic data. (1) UGC number, (2) alternative name, (3) RA, (4) Dec, (5) morphological type, (6) absolute B-band magnitude, (7) heliocentric radial velocity, (8) distance, (9) inclination angle (superscripts indicate source: <sup>t</sup> tilted ring fits to velocity fields, <sup>L</sup> LEDA and <sup>e</sup> estimated from our optical data) and (10) group/cluster membership (superscripts indicate source: <sup>1</sup> Garcia (1993), <sup>2</sup> Tully (1988), <sup>3</sup> Tully (1980), <sup>4</sup> Hickson et al. (1989) and <sup>5</sup> this study). Columns (3) – (5) were taken from NED, column (6) from LEDA and columns (7) and (8) from this study (cf. Table 3).

UGC	alternative	R	A (2	000)	Dec	c (2000)	Туре	M <sub>B</sub>	$V_{hel}$	D	i	group/cluster
	name	h	m	s	0	, <i>ii</i>	31	mag	$km s^{-1}$	Mpc	0	membership
(1)	(2)		(3)	)		(4)	(5)	(6)	(7)	(8)	(9)	(10)
89	NGC 23	0	9	53.4	25	55 26	SB(s)a	-21.48	4555	62.1	50 <sup>t</sup>	LGG 2 <sup>1</sup>
94	NGC 26	0	10	25.9	25	49 55	SA(rs)ab	-20.21	4589	62.6	$42^{t}$	LGG 2 <sup>1</sup>
232	_	0	24	38.7	33	15 22	SB(r)a	-19.96	4839	66.3	$48^L$	_
499	NGC 262	0	48	47.1	31	57 25	SA(s)0/a:	-20.09	4534	62.0	$46^L$	LGG 14 <sup>1</sup>
508	NGC 266	0	49	47.8	32	16 40	SB(rs)ab	-21.97	4647	63.6	$25^{t}$	LGG 14 <sup>1</sup>
624	NGC 338	1	0	36.4	30	40 8	Sab	-21.42	4772	65.1	$59^t$	LGG 14 <sup>1</sup>
798	IC 1654	1	15	11.9	30	11 41	(R)SB(r)a	-19.67	4896	66.7	$40^{t}$	LGG 18 <sup>1</sup>
1310	NGC 694	1	50	58.5	21	59 51	S0? pec	-19.23	2958	40.2	$47^e$	LGG 34 <sup>1</sup>
1541	NGC 797	2	3	27.9	38	7 1	SAB(s)a	-21.10	5649	77.0	$41^{t}$	pair with NGC 801 <sup>5</sup>
2045	NGC 972	2	34	13.4	29	18 41	Sab	-20.39	1527	21.4	$61^e$	$52 - 0^2$ (Cetus)
2141	NGC 1012	2	39	14.9	30	9 6	S0/a?	-18.02	987	14.3	$90^e$	17 -3 <sup>2</sup> (Triang. Spur)
2154	NGC 1023	2	40	24.0	39	3 48	SB(rs)0-	-20.23	695	10.9	$64^L$	N 1023 group <sup>3</sup>
2183	NGC 1056	2	42	48.3	28	34 27	Sa:	-19.60	1540	21.5	$90^e$	_
2487	NGC 1167	3	1	42.4	35	12 21	SA0-	-21.73	4950	67.4	$36^{t}$	LGG 80 <sup>1</sup>
2916	_	4	2	33.8	71	42 21	Sab	-20.34	4518	63.5	$46^{t}$	LGG 115 <sup>1</sup>
2941	IC 357	4	3	44.0	22	9 33	SB(s)ab	-21.38	6261	83.9	$45^e$	group with UGC 2942/435
2953	IC 356	4	7	46.9	69	48 45	SA(s)ab pec	-21.49	894	15.1	$50^{t}$	12 +15 <sup>2</sup> (U. Major)
3205	_	4	56	14.8	30	3 8	Sab	-21.47	3587	48.7	$67^{t}$	_
3354	_	5	47	18.2	56	6 44	Sab:	-20.00	3084	43.6	$90^e$	_
3382	_	5	59	47.7	62	9 29	SB(rs)a	-20.62	4501	62.8	$16^t$	_
3407	_	6	9	8.1	42	5 7	Sa	-20.05	3606	49.8	$45^e$	_
3426	_	6	15	36.3	71	2 15	S0:	-20.89	4005	56.6	$37^e$	LGG 135 <sup>1</sup>
3546	NGC 2273	6	50	8.7	60	50 45	SB(r)a	-20.16	1838	27.3	$52^{t}$	$24 - 1^2$ (Lynx)
3580	_	6	55	30.8	69	33 47	SA(s)a pec:	-18.19	1200	19.2	$63^{t}$	12 -0 <sup>2</sup> (U. Major)
3642	_	7	4	20.3	64	1 13	SA0	-21.00	4498	62.9	$45^{t}$	LGG 140 <sup>1</sup>
3965	IC 2204	7	41	18.1	34	13 56	(R)SB(r)ab	-18.64	4588	62.5	$10^e$	_
3993	_	7	55	44.0	84	55 35	S0?	-19.61	4366	61.9	$20^{t}$	pair with UGC 3992 <sup>5</sup>
4458	NGC 2599	8	32	11.3	22	33 38	SAa	-21.12	4757	64.2	$27^{t}$	pair with PGC 23972 <sup>5</sup>
4605	NGC 2654	8	49	11.9	60	13 16	SBab: sp	-20.08	1350	20.9	$90^e$	13 -6 <sup>2</sup> (U. Major S. Spur)
4637	NGC 2655	8	55	37.7	78	13 23	SAB(s)0/a	-20.97	1415	22.4	$30^{e}$	12 -10 <sup>2</sup> (U. Major)
4666	NGC 2685	8	55	34.7	58	44 4	(R)SB0+ pec	-19.21	876	14.6	$61^L$	13 -4 <sup>2</sup> (U. Major S. Spur)
4862	NGC 2782	9	14	5.1	40	6 49	SAB(rs)a	-20.99	2540	35.9	$20^{e}$	$21 - 0^2$ (Leo)
5060	NGC 2893	9	30	17.0	29	32 24	(R)SB0/a	-18.14	1699	24.1	$24^e$	$21 - 0^2$ (Leo)
5253	NGC 2985	9	50	22.2	72	16 43	(R')SA(rs)ab	-20.72	1325	21.1	$38^t$	12 -7 <sup>2</sup> (U. Major)
5351	NGC 3067	9	58	21.0	32	22 12	SAB(s)ab?	-19.35	1486	21.5	$70^{L}$	21 -12 <sup>2</sup> (Leo)
5559	NGC 3190	10	18	5.6	21	49 55	SA(s)a pec sp	-19.88	1308	18.6	$67^L$	HCG 44 <sup>4</sup>
5906	NGC 3380	10	48	12.2	28	36 07	(R')SBa?	-18.69	1601	23.1	29 <sup>e</sup>	LGG 227 <sup>1</sup>
5960	NGC 3413	10	51	20.7	32	45 59	S0	-17.39	643	10.5	$77^t$	_
6001	NGC 3442	10	53	8.1	33	54 37	Sa?	-17.77	1730	25.1	$47^L$	21 -9 <sup>2</sup> (Leo)
6118	NGC 3504	11	3	11.2	27	58 21	(R)SAB(s)ab	-20.28	1536	22.2	$24^e$	21 -7 <sup>2</sup> (Leo)
6283	NGC 3600	11	15	52.0	41	35 29	Sa?	-17.17	713	12.0	90e	15 -9 <sup>2</sup> (Leo Spur)
6621	NGC 3786	11	39	42.5	31	54 33	(R')SAB(r)a pec	-19.74	2742	38.7	$58^L$	13 -0 <sup>2</sup> (U. Major S. Spur)
6623	NGC 3788	11	39	44.6	31	55 52	SAB(rs)ab pec	-20.24	2670	37.7	$74^L$	13 -0 <sup>2</sup> (U. Major S. Spur)
6742	NGC 3870	11	45	56.6	50	12 00	S0?	-17.28	752	13.0	$36^L$	12 -1 <sup>2</sup> (U. Major)
6786	NGC 3900	11	49	9.4	27	1 19	SA(r)0+	-19.94	1799	25.9	$65^t$	13 -9 <sup>2</sup> (U. Major S. Spur)
6787	NGC 3898	11	49	15.4	56	5 4	SA(s)ab	-20.25	1171	18.9	67 <sup>t</sup>	12 -3 <sup>2</sup> (U. Major)
7166	NGC 4151	12	10	32.6	39	24 21	(R')SAB(rs)ab:	-20.84	998	15.9	$20^t$	12 -6 <sup>2</sup> (U. Major)
7256	NGC 4203	12	15	5.0	33	11 50	SAB0-:	-19.47	1091	16.9	$40^{t}$	14 -1 <sup>2</sup> (Coma–Sculptor)
7489	NGC 4369	12	24	36.2	39	22 59	(R)SA(rs)a	-18.87	1027	16.4	$17^{L}$	12 -6 <sup>2</sup> (U. Major)
7506	NGC 4384	12	25	12.0	54	30 22	Sa	-19.70	2532	37.0	$41^L$	42 +11 <sup>2</sup> (Canes Venatici)
7704	NGC 4509	12	33	6.8	32	5 30	Sab pec?	-16.17	937	14.8	$60^{e}$	LGG 279 <sup>1</sup>
7989	NGC 4725	12	50	26.6	25	30 3	SAB(r)ab pec	-21.69	1208	18.2	$51^{t}$	14 -2 <sup>2</sup> (Coma–Sculptor)
8271	NGC 5014	13	11	31.2	36	16 55	Sa? sp	-18.44	1128	17.7	$60^{e}$	43 -1 <sup>2</sup> (Canes V. Spur)

Table 1. basic data: continued

UGC	alternative	R	A (20	(000	De	c (2000)	Type	$ m M_{B}$	$V_{hel}$	D	i	group/cluster
	name	h	m	S	0	, ,,		mag	$\rm km~s^{-1}$	Mpc	0	membership
(1)	(2)		(3)	)		(4)	(5)	(5) (6) (7) (8) (9)		(10)		
8699	NGC 5289	13	45	8.7	41	30 12	(R)SABab:	-19.70	2521	36.7	$72^{t}$	LGG 361 <sup>1</sup>
8805	NGC 5347	13	53	17.8	33	29 27	(R')SB(rs)ab	-19.69	2384	34.5	$30^{e}$	42 -0 <sup>2</sup> (Canes Venatici)
8863	NGC 5377	13	56	16.7	47	14 8	(R)SB(s)a	-20.32	1791	27.2	$51^{t}$	42 -0 <sup>2</sup> (Canes Venatici)
9133	NGC 5533	14	16	7.7	35	20 38	SA(rs)ab	-21.47	3861	54.3	$53^t$	LGG 380 <sup>1</sup>
9644	_	14	59	34.3	27	6 58	SB(r)a	-19.25	6665	91.5	$19^e$	_
10448	NGC 6186	16	34	25.5	21	32 27	(R')SB(s)a	-21.65	11351	153.8	$30^{e}$	_
11269	NGC 6667	18	30	39.8	67	59 13	SABab? pec	-20.11	2581	38.3	$56^L$	$70 - 0^2 (-)$
11670	NGC 7013	21	3	33.6	29	53 51	SA(r)0/a	-19.53	775	12.7	$68^{t}$	65 +5 <sup>2</sup> (Pegasus Spur)
11852	_	21	55	59.3	27	53 54	SBa?	-20.25	5846	80.0	$50^{t}$	_
11914	NGC 7217	22	7	52.4	31	21 33	(R)SA(r)ab	-20.45	949	14.9	$31^{t}$	65 +1 <sup>2</sup> (Pegasus Spur)
11951	NGC 7231	22	12	30.1	45	19 42	SBa	-19.14	1086	17.4	$70^L$	_
12043	NGC 7286	22	27	50.5	29	5 45	S0/a	-17.15	1007	15.4	$67^{t}$	_
12276	NGC 7440	22	58	32.5	35	48 9	SB(r)a	-20.59	5662	77.7	$37^L$	_
12713	_	23	38	14.4	30	42 29	S0/a	-14.35	295	5.7	$72^{t}$	65 +42 (Pegasus Spur)
12815	NGC 7771	23	51	24.9	20	6 43	SB(s)a	-21.60	4307	58.5	$61^L$	LGG 483 <sup>1</sup>

**Table 2.** Observational parameters: (1) UGC number, (2) observation date, (3) spatial resolution of the unsmoothed data, (4) velocity resolution before Hanning smoothing, (5) rms noise level in the unsmoothed channel maps, (6) 1 sigma column density in the HI maps shown in the atlas, (7) previous 21cm synthesis observations (references are explained at the bottom of the table) and (8) remarks.

UGC	observation date	resolution	$\sigma_{ m ch}$	$\sigma_{ m map}$	previous	remarks
		$RA \times Dec \times v$	CII	тар	observations	
		$^{\prime\prime}$ $^{\prime\prime}$ km s <sup>-1</sup>	mJy/beam	$10^{19} \text{ atoms cm}^{-2}$		
(1)	(2)	(3) (4)	(5)	(6)	(7)	(8)
89	Sep. 2000	$16.1 \times 37.9 \times 10.03$	0.64	3.25	O93	)
94	Sep. 2000	$16.1 \times 37.9 \times 10.03$	0.64	2.86	O93	1 pointing
232	July 2000	$15.7 \times 29.9 \times 10.04$	0.78	4.97	_	,
499	Aug. 2000	$11.1 \times 19.4 \times 10.04$	0.69	2.56	H82, S87	)
508	Aug. 2000	$11.1 \times 19.4 \times 10.04$	0.69	2.82	_	1 pointing
624	July 2000	$10.7 \times 24.2 \times 10.04$	0.61	11.3	_	,
798	Aug. 2000	$12.6 \times 23.2 \times 10.04$	0.73	3.13	_	
1310	June 2000	$12.0 \times 31.8 \times 9.98$	0.90	3.66	VM88	
1541	Aug. 2000	$16.1 \times 27.0 \times 10.07$	0.79	7.06	_	
2045	Aug. 2000	$15.9 \times 34.0 \times 9.94$	0.67	2.75	_	
2141	Jan. 1994 <sup>†</sup>	$8.3 \times 19.4 \times 4.98$	3.51	71.8	_	
2154	July 1997 <sup>†</sup>	$10.1 \times 18.0 \times 19.80$	1.20	2.88	A79, S84	
2183	Apr. 2000	$10.6 \times 26.0 \times 9.94$	0.79	6.52	-	
2487	Dec. 2000	$11.8 \times 22.2 \times 10.06$	0.87	1.29	_	
2916	Apr. 2000	$10.8 \times 13.2 \times 10.03$	0.59	16.8	_	
2941	July 2000	$10.7 \times 32.3 \times 5.06$	1.31	6.47	_	
2953	Sep. 2001	$10.7 \times 32.3 \times 9.00$ $10.7 \times 14.4 \times 9.91$	0.60	1.39	_	
3205	Dec. 2000	$10.6 \times 24.5 \times 10.01$	0.61	15.5	_	
3354	Nov. 2000	$16.9 \times 21.1 \times 9.98$	0.57	5.93	_	
3382	Dec. 2000	$12.2 \times 14.1 \times 5.04$	0.75	1.20	_	
3407	Dec. 2000	$12.2 \times 14.1 \times 3.04$ $12.7 \times 17.8 \times 10.01$	0.73	2.38		
3426	Jan. 2002	$12.7 \times 17.8 \times 10.01$ $12.3 \times 13.0 \times 5.03$	0.71	1.20		
3546	Dec. 2000	$12.3 \times 13.0 \times 9.05$ $12.2 \times 14.9 \times 9.95$	0.63	1.13	VW83	
3580	Jan. 1994 <sup>†</sup>	$9.1 \times 10.2 \times 4.98$	2.50	5.45	B94	
3642	May 2000	$12.6 \times 13.4 \times 10.03$	0.75	3.69	_	
3965	Apr. 2000	$12.0 \times 13.4 \times 10.03$ $11.7 \times 21.4 \times 2.51$	1.84	1.95	_	
3993	July 2000	$12.2 \times 12.3 \times 5.03$	0.82	1.32	_	
4458	July 2002	$12.2 \times 12.5 \times 9.03$ $15.5 \times 43.6 \times 10.04$	0.32	3.14	_	
4605	Apr. 2002	$13.3 \times 43.0 \times 10.04$ $12.8 \times 13.7 \times 9.94$	0.74	18.4	_ W88	
4637	July 2002	$12.6 \times 13.7 \times 9.94$ $12.6 \times 12.4 \times 9.94$	0.68	0.74	<b>VV</b> 00	
4666	Aug. 2000	$12.0 \times 12.4 \times 9.94$ $11.8 \times 15.6 \times 9.91$	0.69	2.95	_ S80	
4862	July 2000	$15.5 \times 24.8 \times 5.00$	0.09	6.07	S94	
5060	July 2000 July 2000	$15.6 \times 33.2 \times 4.99$	0.88	2.04	374	
5253	Aug. 1993 <sup>†</sup>	$13.6 \times 33.2 \times 4.99$ $11.6 \times 12.4 \times 19.92$			O93	
5351	Aug. 1993 Aug. 2000	$11.0 \times 12.4 \times 19.92$ $14.0 \times 20.8 \times 9.94$	1.82 0.75	3.64 4.11	C89, C92	
	Aug. 2000 June 1996 <sup>†</sup>			31.1	W91	
5559 5006		$9.8 \times 26.3 \times 19.92$	1.98	2.04		
5906	Aug. 2000	$34.6 \times 15.8 \times 4.99$	1.05		_	
5960	Apr. 2000	$21.7 \times 11.6 \times 4.97$	1.40	6.45	_	
6001	Aug. 2000	$21.4 \times 11.7 \times 4.99$	0.97	5.07	- VIMO21	
6118	July 2000	$35.2 \times 15.7 \times 9.94$	0.78	4.35	VM83b	
6283	Apr. 1996 <sup>†</sup>	$18.7 \times 12.5 \times 4.97$	3.11	28.9	GR02	)
6621	July 2002	$31.2 \times 15.9 \times 9.97$	0.59	3.70	O93	1 pointing
6623	July 2002	$31.2 \times 15.9 \times 9.97$	0.59	3.70	O93	J 1 2
6742	Apr. 2000	$14.1 \times 12.8 \times 4.97$	0.81	1.84	- -	
6786	Aug. 2000	$15.5 \times 36.0 \times 9.95$	0.74	1.48	VD89, H00	
6787	Apr. 2000	$12.6 \times 13.8 \times 9.92$	0.65	2.14	VD94	
7166	June 1997 <sup>†</sup>	$20.8 \times 12.9 \times 4.98$	2.98	1.73	B77, P92, M99	

<sup>†</sup> observed before the upgrade of the WSRT.

Table 2. observational parameters: continued

UGC	observation date	resolution	$\sigma_{ m ch}$	$\sigma_{ m map}$	previous	remarks
		$RA \times Dec \times v$			observations	
		$^{\prime\prime}$ $^{\prime\prime}$ km s <sup>-1</sup>	mJy/beam	$10^{19} \text{ atoms cm}^{-2}$		
(1)	(2)	(3) (4)	(5)	(6)	(7)	(8)
7256	July 2000	$30.7 \times 15.3 \times 4.98$	1.76	2.29	VD88	
7489	Aug. 1997 <sup>†</sup>	$18.2 \times 11.4 \times 4.98$	4.30	5.18	_	
7506	Aug. 2000	$13.0 \times 10.3 \times 5.00$	0.93	13.1	_	
7704	Aug. 2000	$21.3 \times 10.0 \times 4.98$	0.93	4.69	_	
7989	Oct. 1994 <sup>†</sup>	$8.0 \times 13.2 \times 19.80$	1.13	9.73	W84	
8271	July 2000	$10.2 \times 19.9 \times 4.98$	0.91	0.98	_	
8699	Aug. 2000	$16.9 \times 10.7 \times 9.97$	0.68	12.4	VM83a	
8805	July 1997 <sup>†</sup>	$19.5 \times 10.6 \times 5.00$	4.41	2.49	_	
8863	Aug. 2000	$15.8 \times 22.4 \times 9.95$	0.66	2.11	_	
9133	July 2000	$15.8 \times 28.8 \times 10.02$	0.70	2.44	B92, B94	
9644	July 2000	$10.5 \times 25.6 \times 5.08$	1.08	9.41	_	
10448	July 2000	$10.6 \times 32.1 \times 5.15$	1.10	5.26	_	
11269	July 2000	$12.2 \times 14.4 \times 9.97$	0.59	1.78	_	
11670	May 2000	$11.3 \times 25.0 \times 9.91$	0.73	2.37	K84	
11852	July 2000	$16.1 \times 35.7 \times 10.08$	0.78	4.12	_	
11914	Sep. 2000	$10.5 \times 23.5 \times 9.92$	0.68	2.25	VM95	
11951	Jan. 1994 <sup>†</sup>	$10.1 \times 14.1 \times 4.98$	3.04	54.5	_	
12043	Apr. 2000	$9.9 \times 25.9 \times 4.98$	1.22	9.50	_	
12276	July 2002	$10.4 \times 19.3 \times 5.05$	0.84	9.45	_	
12713	June 2000	$11.6 \times 22.0 \times 4.97$	0.95	1.78	VD91	
12815	Aug. 2000	$10.4 \times 33.4 \times 10.03$	0.70	9.62		

<sup>†</sup> observed before the upgrade of the WSRT.

explanation of the references in column (7):

Allsopp (1979)	S87	Simkin et al. (1987)
Bosma et al. (1977)	S94	Smith (1994)
Broeils (1992)	VD88	van Driel et al. (1988)
Broeils & van Woerden (1994)	VD89	van Driel et al. (1989)
Carilli et al. (1989)	VD91	van Driel & van Woerden (1991)
Carilli & van Gorkom (1992)	VD94	van Driel & van Woerden (1994)
García-Ruiz et al. (2002)	VM83a	van Moorsel (1983a)
Heckman et al. (1982)	VM83b	van Moorsel (1983b)
Haynes et al. (2000)	VM88	van Moorsel (1988)
Knapp et al. (1984)	VM95	Verdes-Montenegro et al. (1995)
Mundell et al. (1999)	VW83	van Woerden et al. (1983)
Oosterloo & Shostak (1993)	W84	Wevers et al. (1984)
Pedlar et al. (1992)	W88	Warmels (1988a)
Shane (1980)	W91	Williams et al. (1991)
Sancisi et al. (1984)		
	Bosma et al. (1977) Broeils (1992) Broeils & van Woerden (1994) Carilli et al. (1989) Carilli & van Gorkom (1992) García-Ruiz et al. (2002) Heckman et al. (1982) Haynes et al. (2000) Knapp et al. (1984) Mundell et al. (1999) Oosterloo & Shostak (1993) Pedlar et al. (1992) Shane (1980)	Bosma et al. (1977) S94 Broeils (1992) VD88 Broeils & van Woerden (1994) VD89 Carilli et al. (1989) VD91 Carilli & van Gorkom (1992) VD94 García-Ruiz et al. (2002) VM83a Heckman et al. (1982) VM83b Haynes et al. (2000) VM88 Knapp et al. (1984) VM95 Mundell et al. (1999) VW83 Oosterloo & Shostak (1993) W84 Pedlar et al. (1992) W88 Shane (1980) W91

**Table 3.** Hi properties: (1) UGC number, (2) systemic velocity, (3) distance, (4) line width of global profile at the 20% level, (5) idem at the 50% level, (6) total Hi flux, (7) total Hi mass, (8) Hi radius in ", (9) idem in kpc, (10) average face-on Hi surface density within Hi radius, (11) idem within optical radius, (12) amount of asymmetry in global profile, (13) morphological asymmetry, (14) kinematical asymmetry, and (15) evidence for current interaction/merging/accretion.

UGC	$V_{ m sys}$	D	$W_{20}^c$	$W_{50}^c$	∫ F dv	$M_{ m HI}$	R <sub>HI</sub>	R <sub>HI</sub>	$<\Sigma_{\rm HI}>_{\rm R_{ m HI}}$	$<\Sigma_{\rm HI}>_{\rm R_{25}}$		Asymi	metries	
	km s <sup>-1</sup>	Mpc	$km s^{-1}$	$km s^{-1}$	$Jy \text{ km s}^{-1}$	$10^9 \mathrm{M}_{\odot}$	"	kpc	$M_{\odot}  \mathrm{pc}^{-2}$	$M_{\odot} pc^{-2}$	prof.	morph.	kin.	interac.
(1)	(2)	(3)	(4)	(5)	(6)	(7)	(8)	(9)	(10)	(11)	(12)	(13)	(14)	(15)
89	4555	62.1	440	386	9.44	8.63	77	24	3.5	5.1	*	*		+
94	4589	62.6	319	300	12.19	11.28	85	26	4.0	7.2	*	*		+
232	4839	66.3	276	236	7.66	7.95	82	27	2.1	3.3		*		
499	4534	62.0	91	65	23.9	21.45	104	31	1.6	1.8		**	n.a.	+
508	4647	63.6	478	453	7.71	7.40	94	29	2.0	2.0		*		+
624	4772	65.1	560	518	14.39	14.42	92	29	4.3	7.7		**	***	+
798	4896	66.7	220	205	3.88	4.08	72	23	1.8	2.4				
1310	2958	40.2	186	117	6.16	2.33	50	9.7	4.8	14.3	**	n.a.	n.a.	+
1541	5649	77.0	446	421	7.94	11.14	75	28	2.8	3.9	***	*	*	
2045	1527	21.4	334	289	18.89	2.09	94	9.9	4.2	3.6				
2141	987	14.3	232	201	47.94	2.30	155	10	4.5	11.3		*		
2154	695	10.9	529	497	80.13	2.26	_#	_#	_#	0.7	n.a.	n.a.	n.a.	+
2183	1540	21.5	292	268	32.99	3.63	164	17	3.0	6.8	*	*	*	
2487	4950	67.4	466	437	15.93	17.09	120	39	1.6	1.8				
2916	4518	63.5	357	337	23.08	21.94	108	34	4.3	7.3	*		**	+
2941	6261	83.9	160	139	9.84	16.32	75	30	3.4	5.4		*	*	+
2953	894	15.1	484	464	119.73	6.43	300	22	1.9	2.2				
3205	3587	48.7	437	419	16.38	9.18	118	28	3.4	4.7				
3354	3084	43.6	423	398	17.71	7.95	77	16	2.0	2.3	*			
3382	4501	62.8	204	193	5.68	5.27	78	24	2.0	2.8				
3407	3606	49.8	316	298	3.13	1.82	50	12	2.6	3.1	*			+
3426	4005	56.6	161	114	3.27	2.53	_#	_#	_#	0.2	*	n.a.	n.a.	+
3546	1838	27.3	367	347	13.74	2.42	132	18	2.0	2.2				
3580	1200	19.2	241	224	43.24	3.76	199	18	3.0	5.4	*		*	
3642	4498	62.9	469	441	38.13	35.62	171	52	3.1	4.7	*			
3965	4588	62.5	93	74	15.67	14.47	106	33	3.2	5.3				
3993	4366	61.9	209	192	7.73	6.99	82	24	1.6	1.7		*		
4458	4757	64.2	284	245	12.24	11.82	110	34	2.0	3.1		**		+
4605	1350	20.9	432	392	58.03	5.99	277	28	1.8	3.0				
4637	1415	22.4	399	243	22.34	2.62	_#	_#	_#	0.9	**	***	n.a.	+
4666	876	14.6	303	284	29.07	1.47	184	12	2.2	2.5	*	*		
4862	2540	35.9	192	156	16.42	5.07	90	16	3.1	2.2	***	***	n.a.	+
5060	1699	24.1	185	86	5.13	0.70	60	7.0	2.3	3.7	**	***		
5253	1325	21.1	320	296	132.4	13.92	330	34	2.1	3.8	*	***		+
5351	1486	21.5	276	187	9.23	0.99	61	6.3	3.4	3.1	**	**	**	+
5559	1308	18.6	473	449	7.32	0.60	88	7.7	2.6	1.5		n.a.	n.a.	+
5906	1601	23.1	123	110	2.53	0.31	47	5.3	2.0	1.9		*		
5960	643	10.5	180	154	28.42	0.75	109	5.7	6.4	15.9	**	*	**	
6001	1730	25.1	153	121	3.58	0.33	44	5.3	4.2	9.6	*	n.a.	n.a.	
6118	1536	22.2	219	194	7.20	0.84	89	9.5	2.7	3.0	*			
6283	713	12.0	216	200	57.38	2.00	204	12	3.1	5.5		*		
6621	2742	38.7	420	375	3.69†	$1.30^{\dagger}$	57	10	4.1	4.1		***	n.a.	+
6623	2670	37.7	520	438	$11.51^{\dagger}$	$3.95^{\dagger}$	79	15	3.5	4.4	***	***	n.a.	+
6742	752	13.0	127	84	4.08	0.17	51	3.2	4.3	8.0	*	**		
6786	1799	25.9	445	423	24.91	3.93	143	18	2.4	3.4				
6787	1171	18.9	484	466	46.82	3.96	256	23	1.7	2.5	*	*		+
7166	998	15.9	140	121	67.58	4.05	304	24	2.1	2.2	**	*	*	
7256	1091	16.9	271	239	48.11	3.21	226	19	1.3	1.8	*	**	n.a.	+
7489	1027	16.4	86	68	4.07	0.26	55	4.4	3.4	2.7		n.a.	n.a.	
7506	2532	37.0	171	116	3.66	1.17	37	6.7	4.1	3.7	*	n.a.	n.a.	
7704	937	14.8	93	61	6.49	0.34	58	4.3	3.7	9.4		*		
7989	1208	18.2	417	390	125.78	9.78	380	34	2.4	2.7	***	**	*	+
8271	1128	17.7	166	76	10.37	0.78	89	7.7	2.1	3.7	*	***	n.a.	+

<sup>&</sup>lt;sup>#</sup> Average surface density  $< 1 \, \rm M_{\odot} \, pc^{-2}$  everywhere.

<sup>†</sup> Total flux and H<sub>I</sub> mass measured from 30" data.

Table 3. HI properties: continued

UGC	$V_{ m sys}$	D	$W_{20}^{c}$	$W_{50}^{c}$	∫ F dv	$M_{ m HI}$	R <sub>HI</sub>	R <sub>HI</sub>	$<\Sigma_{\rm HI}>_{\rm R_{\rm HI}}$	$<\Sigma_{\rm HI}>_{\rm R_{25}}$	Asymmetries			
	$km s^{-1}$	Mpc	$km s^{-1}$	$\mathrm{km}\mathrm{s}^{-1}$	Jy km s <sup>-1</sup>	$10^9 \mathrm{M}_{\odot}$	"	kpc	$M_{\odot} pc^{-2}$	$M_{\odot} pc^{-2}$	prof.	morph.	kin.	interac.
(1)	(2)	(3)	(4)	(5)	(6)	(7)	(8)	(9)	(10)	(11)	(12)	(13)	(14)	(15)
8699	2521	36.7	387	368	10.22	3.26	100	17	2.9	4.6				
8805	2384	34.5	129	82	15.45	4.34	137	23	1.8	3.0				
8863	1791	27.2	391	367	14.44	2.52	148	19	1.5	1.6	*		*	
9133	3861	54.3	451	416	43.25	30.23	218	57	1.9	4.0	**	**		
9644	6665	91.5	131	119	3.63	7.16	53	24	3.2	4.4				
10448	11351	153.8	99	85	1.76	9.82	40	29	2.6	2.0				
11269	2581	38.3	415	344	31.32	10.86	161	30	3.0	6.1	*	***	**	+
11670	775	12.7	342	325	28.05	1.08	158	9.9	2.5	2.5	**		*	
11852	5846	80.0	329	305	19.55	29.6	137	53	2.5	4.2	**		**	
11914	949	14.9	322	309	13.16	0.69	113	8.1	3.0	2.8				
11951	1086	17.4	222	201	25.93	1.85	127	10.7	4.4	9.6				+
12043	1007	15.4	190	176	20.96	1.18	143	10.8	2.4	6.1				
12276	5662	77.7	135	108	4.09	5.83	67	26	2.2	2.9	*	*		+
12713	295	5.7	147	101	10.79	0.08	84	2.3	2.0	4.3	*	*	*	
12815	4307	58.5	656	457	11.74\$	9.36\$	103	30	1.9	2.6	n.a.	n.a.	n.a.	+

<sup>\$</sup> Includes emission from UGC 12813.

# Appendix A: Atlas of H<sub>I</sub> observations

On the following pages, we present the H<sub>I</sub> observations for all our sample galaxies. For each galaxy, we show a figure consisting of 6 panels:

**Upper left:** Grayscale and contour image of the integrated H<sub>I</sub> distribution. Thin contours are at 2, 4, 8, 16, ....  $\sigma_{map}$ , with  $\sigma_{map}$  determined as in Sect. 3.4; its value in atoms cm<sup>-2</sup> is given in Table 2. The thick contour indicates the 1  $M_{\odot}$  pc<sup>-2</sup> level, corrected to face-on. For edgeon galaxies, the correction to face-on could not be done, and we only show the  $\sigma_{map}$ -levels. The beam size for the data shown here is given in the lower left. Note that this is not always the same as the resolution given in Table 2, as we sometimes needed to smooth the data in order to increase the signal-to-noise ratio.

**Lower left:** The same image overlayed on an optical R-band image. We show only the  $2, 8, 32, 128, \ldots \sigma_{map}$  and  $1 \text{ M}_{\odot} \text{ pc}^{-2}$  contours here, and omit the contours at  $4, 16, 64, \ldots \sigma_{map}$ . The images were taken as part of a parallel project to image our sample galaxies at optical wavelengths, and will be described elsewhere (Noordermeer et al., in prep.).

**Upper middle:** Grayscale and contour representation of the velocity field. Dark shading indicates the receding side. Contours are spaced at 25 km s<sup>-1</sup>. The thick contour indicates the systemic velocity from Table 3.

**Lower middle:** Position-velocity diagram along the major axis. Contours and grayscales show a slice through the data cube, the squares show the corresponding velocities from the velocity field, determined by fitting skewed Gaussians to the line profiles (cf. Sect. 3.5). Note that these are *not* rotation curves, and that in some cases there are clear and pronounced deviations from the projected rotational velocities (e.g. UGC 2183; cf. Sect. 3.5 and individual notes in Sect. 5). Contours are at -1.5 and -3 (dotted) and 1.5, 3, 6, 12, ... times the rms noise in the channel maps. The position angle along which the slice is taken is indicated in the top left, the systemic velocity from Table 3 and central position are indicated with the dashed lines.

**Upper right:** Global profile. The dashed line indicates the systemic velocity from Table 3. Note that we always used the Hanning-smoothed data cube at 60" resolution to produce the global profile (except for UGC 6621 and 6623, see Sect. 5).

**Lower right:** The radial H<sub>I</sub> surface density profile. The dotted line indicates the profile for the receding half of the galaxy, the dashed line the approaching side. The thick solid line gives the profile averaged over the entire disk. Data are only shown when above the  $2\sigma_{\text{map}}$ -level in the H<sub>I</sub>-map. Profiles with symbols are determined by measuring the intensity on concentric ellipses; profiles without symbols are for edge-on galaxies and are determined using Warmels's method (see Sect. 3.6 for details).

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